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# NAVAL POSTGRADUATE SCHOOL Monterey, California



# THESIS

"HEAD-ON" SCATTERING OF A TUBULAR CYLINDER
OF FINITE LENGTH
FOR RADAR TARGET IDENTIFICATION PURPOSES

bν

David Geller

March 1985

Thesis Advisor:

Hung-Mou-Lee

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the investigation of the cross section of a target over a broad frequency band.

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"Head-on" Scattering of a Tubular Cylinder of Finite Length for Radar Target Identification Purposes.

bу

David Geller Lieutenant, Išraeli Navy B.S., Technion-Israeli High Institute of Technology,1977

Submitted in partial fulfillment of the requirements for the degree of

MASTER OF SCIENCE IN ELECTRICAL ENGINEERING

from the

NAVAL POSTGRADUATE SCHOOL March 1985

#### ABSTRACT

This thesis studies the "Head-on" back scattering of a finite tubular cylinder with a circular cross-section and a very thin conducting wall.

At this aspect angle,the back scattered fields depend only on the first Fourier component of the circumferential variations of the  $\varphi$  -current.

Measurements of several scaled tubular cylinders were taken and the experimental results were compared to theoretical data available.

This thesis is part of an ongoing project of target identification through the investigation of the cross-section of a target over a broad frequency band.

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#### I. INTRODUCTION

Identifying a target by its back scattering signal is desirable if different actions are to be taken towards different targets. At a time when warfare is no longer conducted face to face, but missiles have the ability to destroy thier targets long before they come within visual range, it has become necessary to identify targets by some means other than vision. The identification is needed because of the fact that current operational policy requires a positive identification of a target before destroying it. This in a sense completly negates the prime capability of sophisticated weapon-systems. For example, the HARPOON missile whose range is far beyond the horizon prevents a visual identification of the target from the fireing platform. A fact that limits the use of it to only specific situations under sever restrictions.

At present, most identifications are done with radars with which the presence of a target can be detected. In addition to the fact that the presence of a target can be discovered, information about its position, speed and acceleration may be obtained. An experienced radar operator can sometimes distinguish a big target from a small target by its spot-size on the radar screen, but this is not enough to identify the target.

To enhance our capability of identifying a target, it is advantageous to look at resonances exited by electromagnetic fields of different frequencies incident on the target. The reasons for looking at this particular range of frequencies are the following: First, at resonance the scattered fields are stronger compared to non-resonance situations. The

back-scattering cross-sections are larger and the target can be detected more easily. Second, the resonance frequencies and the amplitudes and phase shifts of the scattered fields at these frequencies are determined by the geometry of the target. Since the number of targets of our interest are finite, the identity of a target is revealed by examing a small data base.

This thesis is part of an ongoing project of target-identification through the investigation of the wide-band cross-section of a target. A finite tubular cylinder which has a circular cross section and a very thin conducting wall serves as the canonical target. The finite tubular cylinder is chosen because of its resemblance to a missile body and because theoretical formulations are available for its surface current distribution and scattered fields.

The tubular cylinder was placed in an anechoic chamber and was irradiated with an incident electromagnetic wave. While irradiated by the incident electromagnetic wave. surface currents were excited and generated scattered-field. The surface current had an axial and a circumferential component. The circumferential current circled around the cylinder while the axial current traveled along the cylinder and was reflected at the ends. At the same time, the incident wave kept impinging on the cylinder and excited new surface currents which added to the existing ones. At a certain frequency, the newly excited axial current might add constructively to the current reflected from one end, resulted in a large axial current, which gave a strong scattered field. Or the newly excited circumferential current might add constructively to the current which had made a complete circle around the cylinder, and a strong scattered field could be observed.

This thesis studied the "head-on" back scattering of the cylinder. At this aspect angle, the back scattered fields depended only on the first Fourier component of the circumferential variation of the  $\varphi$ -current. Measurements of several scaled tubular cylinders were taken and the experimental results were compared to theoretical data available.

For the next step of this project, more complicated models evolving from a tubular cylinder to a missile will be constructed and studied. Fins will be added, one end of the cylinder will be closed and rounded to perturb the model further and finally wings will be added to make a true missile model. By comparing the scattering data of the models to those of the tubular cylinders, the effects of the successive perturbations to the physical structure on the back-scattering cross-section and phase shift will be investigated.

Chapter II deals in its first part with electromagnetic back scattering theory in general, the definition of radar cross section , the polarization matrix and methods to obtain it. Its second part contains the solution to the electromagnetic back-scattering of a tubular cylinder with finite length. Chapter III describes some CW step frequency cross section measurements which are carried out, including the experimental setup, the measurement procedures and the measured results. Chapter  $^{\rm IV}$  deals with data analysis, compares the experimental results to theoretical data and presents some conclusions and recommendations for the future.

#### II. ELECTROMAGNETIC SCATTERING THEORY

A typical problem in electromagnetic scattering consists of these main elements.

- 1. Transmitting system: RF source and antennas.
- 2.An object of arbitrary shape and size as a target.
- 3.Receiving system: Antenna and receiving equipment to determine the amplitude, phase and polarization of fields at any point in space.

The transmitting system causes incident fields  $\overset{\cdot}{E}^1$ ,  $\overset{\cdot}{H}^1$  to imping upon the target. The current in the source induces time-varying distributions of oscillating charges and currents in the scatterer. These currents cause scattered, or reradiated fields  $\overset{\cdot}{E}^3$ ,  $\overset{\cdot}{H}^3$ . The total fields  $\overset{\cdot}{E}^3$ ,  $\overset{\cdot}{H}^3$  are the vector sum of the incident and scattered fields.

In order to simplify the theoretical problem, it is usually assumed that the source is not coupled to the target. This fact enables one to obtain the scattered field by subtracting the incident fields from the total fields.

A configuration of the scattering problem is shown in Figure  $2.1\,$ 

This chapter deals with the analytical background, the definition of radar cross-section, the polarization scattering matrix, and finally the specific problem of the scattering by a finite tubular cylinder.

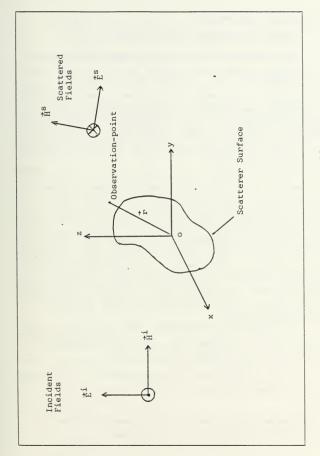


Figure 2.1 The Scattering Problem

#### A. ANALYTICAL BACKGROUND

#### 1. Definition of Radar Cross-section

The radar cross-section of a target is a quantitative measure of the ratio of power density in the vector signal scattered in the direction of the receiver to the power density of the radar wave incident upon the target. The vectorial nature of the electromagnetic interaction requires specifications of the polarization of the incident wave with reference to target orientation in three dimensions. Radar operating frequency is an additional parameter which must be specified.

Thus, a single number specification for the radar cross section holds for a particular target, a specific polarization and the frequency of the incident wave, the aspect angle of the target relative to the incident wave, and the polarization of the receiving antenna.

The radar cross-section is defined to be independent of range to the target. This definition holds under far field assumption. Which means that the target is sufficiently far from the transmitting antenna to justify the assumption that the incident wave is planar at the target and the scattered wave is also planar in the neighborhood of the receiving antenna.

The theoretical definition of the radar cross-section relates incident to scattered electromagnetic fields, as given in equation  $2.1\,$ 

$$\sigma = 4\pi R^2 \lim_{R \to \infty} \left| \frac{E_s}{E_o} \right|^2$$
 (eqn 2.1)

where:

 $\rm E_{\rm O}^{=}$  magnitude of electric field component of incident electromagnetic field at the target.

- ${\bf E}_{\rm S}$  = magnitude of electric field component of scattered electromagnetic field as measured by a hypothetical observer.
- R = distance from target to the hypothetical observer.

The radar cross-section  $\sigma$  has the dimensions of area  $\,m^{\,2}\,.$ 

The limiting process is introduced in equation 2.1 to assure that the distance at which the hypothetical observation is made, is far enough from the target. Under the free-space conditions assumed, the quantity  $\left|E_{\rm S}/E_{\rm O}\right|^2$  is proportional to the power flux density of the scattered waves.

# 2. Polarization Scattering Matrix

The radar cross-section of a target depends upon the target shape and material, the angle (or angles, in a case of bistatic system) at which the target is viewed, radar frequency, and the polarization of the radar-transmitting and receiving antennas.

In particular, if a target is viewed at a specific aspect angle with a single frequency, the radar cross section depends upon polarization [Ref. 1].

The polarization scattering matrix is introduced in order to express target reradiation independent of radar polarization [Ref. 2].

Scattering is expressed as an explicit function of radar polarization, when matrices describing the polarization properties of antennas and target are defined.

The transmitting and receiving antennas can be represented by the matrices:

$$\hat{q} = \begin{bmatrix} \cos\phi_t & - \\ \sin\phi_t e^{j\delta t} \\ - \end{bmatrix}$$
 (eqn 2.2)

$$\hat{p} = [\cos\phi_r \quad \sin\phi_r e^{j\delta r}]$$
 (eqn 2.3)

where:

- q = column matrix defining the polarization of the transmitting antenna.
- $\hat{p}$  = row matrix defining the polarization of the receiving antenna.
- $\phi$  = an angle which when  $\delta$ =0, denotes the orientation of the linear polarization of the antenna when used for transmitting , refered to the horizontal plane.
- $\delta$  = phase angle.
- t = denotes transmitting antenna.
- r = denotes receiving antenna.

The configuration of the cylinder for the measurement in this thesis and the antenna polarization angle are shown in Figure  $2.2\,$ 

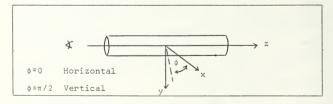


Figure 2.2 The Antenna Polarization Angle

The radar cross-section of a target observed by a transmitting antenna with polarization q and a receiving antenna with polarization p is given by equation 2.4

$$\sigma = |\hat{p}\hat{Sq}|^2 \qquad (eqn 2.4)$$

where S denotes the complex scattering matrix used to represent the polarization properties of the target. Assuming a uniform plane incident wave, the scattering matrix is a linear relation between the incident field and the scattered far field from the target.

With p and q defind by equations 2.2 and 2.3 the scattering matrix S is a  $2 \!\!\!\!\!\!\!^{\pm} 2$  matrix as shown below in equation 2.5

$$S = \begin{bmatrix} \sqrt{\sigma_{HH}} & 1^{j\rho_{HH}} & \sqrt{\sigma_{HV}} & 1^{j\rho_{HV}} \\ \sqrt{\sigma_{VH}} & 1^{j\rho_{VH}} & \sqrt{\sigma_{VP}} & 1^{j\rho_{VV}} \end{bmatrix}$$
 (eqn 2.5)

where:

 $\sqrt{\sigma}$  = magnitude of the scattering matrix element.

ρ = phase of the scattering matrix element.

H = denotes horizontal polarization.

V = denotes vertical polarization.

The scattering matrix is symmetrical  $(\sqrt{\sigma_{H}\bar{y}}\sqrt{\sigma_{VH}}\rho_{HV}^{-a}\rho_{VH})$  in at list two cases .

 Bistatic scattering when the body is a perfect conductor. 2. Back-scattering from an arbitrary body.

The transmitting and receiving antennas used in our case are linearly horizontally polarized.

The expressions for the transmitting and receiving antennas under this condition are:

$$\hat{q} = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$$
 (eqn 2.6)

$$\hat{p} = [1 \ 0]$$
 (eqn 2.7)

since  $\phi = 0$ .

The radar cross-section is:

$$\sigma = \sigma_{HH}$$
 (eqn 2.8)

#### 3. Methods of Obtaining the Scattering Matrix

Ideally one would compute the radar cross-section of a target through the formal solution of Maxwells equations. Those equations should be solved for the boundary conditions appropriate to the target.

 $\label{eq:theorem} \mbox{The major mathematical method for obtaining an exact} \\ \mbox{solution is the separation of variables.}$ 

Formal solutions via separation of variables are possible only for a few special cases. In those cases the wave equation is separable in a coordinate system, having a

coordinate surface that coincides with the surface of the body [Ref. 3]. This situation explains why exact solutions are rare, and for many practical problems, the use of approximations is the only practical approach.

The integral equation formulation shows that electromagnetic scattering of an incident plane wave by an arbitrary body can be described in terms of an integral of various vector products. Those vector products involve the surface electric and magnetic fields. One form which is convenient for this purpose is the Chu-Stratton integral. [Ref. 4]. This integral is an exact representation of the scattered electromagnetic field. It is given in terms of an integration over a complete surface enclosing the body in question. In particular, if there were available knowledge of the total distribution of electric and magnetic fields about the body, insertion of these values in the Chu-Straton integral would permit the immediate solution of the scattering problem.

Numerical schemes have been designed to solve the integral equations approximately. High speed digital computers are needed to establish surface currents flowing on the target. For example surface current distribution can be computed by a finite difference solution to a network of simultaneous equations [Ref. 5].

Because of limitation of computation time and storage capability, the finite difference solution is applicable only when the dimensions of the target do not exceed a very few wavelengths.

For targets larger than a few wavelengths in dimension, asymptotic methods are frequently used. There are three levels of complexity:

The simplest approach is the geometric optics approach. It treats ray bundles by the laws of reflection and refraction [Ref. 6]. However, geometric-optics failes

to distinguish the effects of polarization and the wave nature of the problem.

The second approach is the physical optics. In this approach the local current density, at each point on the illuminated portion of the body, is assumed to be identical to that which would flow at that point on an infinite tangent plane [Ref. 7]. Physical optics is not valid for applications entailing accurate specification effects [Ref. 8].

The third approach is the geometric diffraction theory. This approach is an extension of geometric optic that accounts for diffraction [Ref. 9]. The technique combines the simplicity inherent in the ray approach with the necessary consideration of wavelengths and phases of the wave. But it never include resonances of the target.

#### B. SCATTERING BY A FINITE TUBULAR CYLINDER

When a cylinder is irradiated with an incident electromagnetic wave, surface current is exited on the cylinder. This surface current will radiate and generate the scattered field.

Electromagnetic scattering from conducting objects in a homogeneous, isotropic medium can be treated as a boundary value problem. By use of the Stratton-Chu equations, integrodifferential equations can be set up for the current distribution on the surfaces of the objects, with the Greens function in the medium as the kernel.

The current distribution on the surface of a tubular cylindrical conductor with negligible wall thickness, excited by an incident electromagnetic field, can be written as a pair of coupled integrodifferential equations with the sum of the inside and outside surface currents as the unknown and the incident tangential electric field on

the surface of the conductor as the given quantity. Such a coupled integrodifferential equations were given by Lee [Ref. 16].

$$(1 + \frac{1}{1_1^2}, \frac{9^2}{92^2}) \int_{-1}^{1} dz_0 K_{z,n}(z_0) G_n(1_1 | z-z_0|, 1_2)$$
 (eqn 2.9)

$$+ \ \frac{\mathrm{i}\,n}{1_{1} 1_{2}} \ \frac{\partial}{\partial z} \ \int\limits_{-1}^{1} \ \mathrm{d}z_{o} \ K_{\phi\,n}(z_{o}) G_{n}(1_{1}|z-z_{o}|,1_{2}) = - \ \frac{2\mathrm{i}}{1_{1} 1_{2} \xi_{o}} E_{zn}^{sc}(z) \,.$$

$$\int_{-1}^{1} dz_{o} K_{\phi n}(z_{o}) \left\{ \frac{1}{2} [G_{n-1}(1_{1} | z-z_{o} | , 1_{2}) \right\}$$
 (eqn 2.10)

+ 
$$G_{n+1}(1_1|z-z_0|,1_2)] - \frac{n^2}{1_2^2} G_n(1_1|z-z_0|,1_2)$$

$$+ \frac{in}{l_1 l_2} \frac{\partial}{\partial z} \int_{-1}^{1} dz_0 K_{zn}(z_0) G_n(l_1|z-z_0|, l_2) = - \frac{2i}{l_1 l_2 \xi_0} E_{\phi n}(z).$$

equations 2.9 and 2.10 can be obtained from the Stratton-Chu equations, together with the edge conditions that:

$$K_{z}(0,z) = 0 \cdot (1-z^{2})^{1/2}$$
 as  $|z| \to 1$  (eqn 2.11)

In the previous equations:

sc-denotes-scattered.

i -denotes-incident.

 $\xi = (\mu/\epsilon)^{1/2}$ 

and

$$G_{n}(1_{1}|z-z_{o}|, 1_{2}) = (eqn 2.12)$$

$$\int_{-\pi}^{\pi} \frac{d\phi}{2\pi} 1^{-in(\phi-\phi_{o})} \cdot G[1_{1}|z-z_{o}|, 21_{2}|sin(\phi-\phi_{o})/2|]$$

$$G(x_{1},x_{2}) = (exp[i(x_{1}^{2}+x_{2}^{2})^{-1/2}])(x_{1}^{2}+x_{2}^{2})^{-1/2} (eqn 2.13)$$

The following equations, 2.14 and 2.15 are boundary conditions for the tangential electric field components on a perfectly conducting surface.

$$E_{zn}^{sc}(z) + E_{zn}^{i}(z) = 0$$
 -1

$$E_{\phi n}^{SC}(z) + E_{\phi n}^{i}(z) = 0$$
 -1

The cylinder is assumed to be in a medium with homogeneous isotropic permittivity  $^\epsilon$  and permeability  $^\mu$ . The half length of the cylinder  $1_1$  and the radius  $1_2$  are measured in 1/k ,with k= $\omega\left(\epsilon\mu\right)^{1/2}$  being the wave number, so that  $1_1$ =kh and  $1_7$ ka.

The coordinate system is scaled so that the tubular cylinder occupies the region -1<z<1, $\rho$ =1 as shown in Figure 2.3

On the surface,  $\rho=1$ , there are scattered electric field  $E^{S'}(\phi,Z)$ , the incident electric field  $E^{\dot{1}}(\phi,Z)$  and for -1<Z<1, the surface current  $\vec{K}(\phi,Z)$ , which is the sum of the

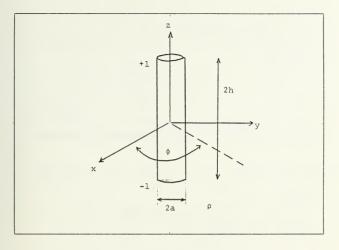


Figure 2.3 The Tubular Cylinder

outer surface current  $\vec{K}^+(\varphi,Z)$  on  $\rho=1^+$  and the inner surface current  $K^-(\varphi,Z)$  on  $\rho=1^-$ .

The surface currents can be represented as:

(eqn 2.16)

$$K_{z}(\phi,z) = \sum_{n=-\infty}^{\infty} e^{jn\phi} K_{zn}(z) = \sum_{n=0}^{\infty} [K_{zn}^{(+)}(z) \cos \phi + iK_{zn}^{(-)}(z) \sin \phi].$$

where

$$K_{ZO}^{(+)}(z) = K_{ZO}(z).$$

$$K_{2n}^{(+)}(z) = K_{2n}(z) + K_{2,-n}(z)$$
  
 $K_{2,0}^{(-)}(z) = 0$ 

$$\mathsf{K}_{\varphi}(\phi\,,z) = \sum_{\mathsf{n}=-\infty}^{\infty} = \phi^{\mathsf{j}\,\mathsf{n}\varphi} \mathsf{K}\phi_{\mathsf{n}}(z) = \sum_{\mathsf{n}=0}^{\infty} \left[\mathsf{K}_{\varphi}^{(+)},\mathsf{n}(z)\mathsf{cosn}\phi + \mathsf{i}\mathsf{K}_{\varphi}^{(-)},\mathsf{n}(z)\mathsf{sinn}\phi\right].$$

where

$$K_{\phi,0}^{(+)}(z) = K_{\phi,0}(z).$$

$$K_{\phi}^{\left(+\right)}(z) = K_{\phi}, n(z) + K_{\phi}, -n(z).$$

$$K_{\phi,0}^{(-)}(z) = 0.$$

 $K_{Z}^{}(\varphi,z)$  is the axial current density,and  $K_{\varphi}^{}(\varphi,z)$  is the circumferential current density.

In the far field where:

$$K|\hat{r}| = (1_2^2 \rho^2 + 1_1^2 z^2)^{1/2} >> 1$$
 (eqn 2.18)

and

$$|K|\hat{r}| > (1_1^2 + 1_2^2)^{1/2} \ge |K|\hat{r}_0|$$
 (eqn 2.19)

$$-\frac{2i}{l_{1}l_{2}\xi_{o}} E_{z}(\rho,\phi,z) = \sin^{2}\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(\overset{+}{r}-\overset{+}{r}_{o}) K_{z}(\phi_{o},z_{o}) - \sin^{2}\theta_{o} G(\overset{+}{r}-\overset{+}{r}_{o}) \sin(\phi-\phi_{o})$$

$$\sin^{2}\theta_{o} G(\overset{+}{r}-\overset{+}{r}_{o}) \sin(\phi-\phi_{o})$$

$$K_{\phi}(\phi_{o},z_{o})$$
(eqn 2.20)

$$\frac{-2i}{1_1 l_2 \tilde{z}_o} E_{\rho}(\rho, \phi, z) = -\sin\theta \cos\theta \int_{-1}^{1} dz_o \int_{-\pi}^{\pi} \frac{d\phi_o}{2\pi} G(\vec{r} - \vec{r}_o) Kz(\phi_o, z_o) + \cos^2\theta \int_{-\pi}^{1} dz_o \int_{-\pi}^{\pi} \frac{d\phi_o}{2\pi} G(\vec{r} - \vec{r}_o) \sin(\phi - \phi_o) K_{\phi}(\phi_o, z_o).$$

$$\frac{-2i}{l_1 l_2 \xi_0} E_{\phi}(\rho,\phi,z) = \int_{-1}^{1} dz_0 \int_{-\pi}^{\pi} \frac{d\phi_0}{2\pi} G(r - r_0) \cos(\phi - \phi_0) K_{\phi}(\phi_0,z_0)$$

(eqn 2.22)

In spherical coordinates, the expressions are simpler. Because:

$$E_{r}(r,\theta,\phi)=E_{0}(\rho,\phi,z)\sin\theta+Ez(\rho,\phi,z)\cos\theta$$
 (eqn 2.23)

$$E_{\theta}(r,\theta,\phi)=E_{\rho}(\rho,\phi,z)\cos\theta-E_{z}(\rho,\phi,z)\sin\theta \tag{eqn 2.24}$$

equations 2.20,2.21 and 2.22 can be written in the following form:

$$-\frac{2i}{l_1 l_2 \xi_0} E_{\theta}(r, \theta, \phi) = 0$$
 (eqn 2.25)

$$-\frac{2i}{l_1 l_2 \xi_0} E_{\theta}(\mathbf{r}, \theta, \phi) = -\sin \theta \int_{-1}^{1} d\mathbf{z}_0 \int_{-\pi}^{\pi} \frac{d\phi_0}{2\pi} G(\mathbf{r} - \mathbf{r}_0) K_z(\phi_0, \mathbf{z}_0) + (eqn 2.26)$$

$$\cos\theta \int_{-1}^{1} dz_{o} \int_{-\pi}^{\pi} \frac{d\phi_{o}}{2\pi} G(\mathring{r} - \mathring{r}_{o}) \sin(\phi - \phi_{o}) K_{\phi}(\phi_{o}, z_{o})$$

$$-\frac{2i}{l_1 l_2 \xi_0} E_{\phi}(\mathbf{r},\theta,\phi) = \int_{-1}^{1} dz_{e_{-\pi}} \int_{-\pi}^{\pi} \frac{d\phi_{e_{-\pi}}}{2\pi} G(\mathring{\mathbf{r}}-\mathring{\mathbf{r}}_{o}) \cos(\phi-\phi_{o}) K_{\phi}(\phi_{o},z_{o})$$

(eqn 2.27)

where

$$\begin{array}{c} -\mathrm{il}_2 \sin\theta \cos(\phi - \phi_0) & -\mathrm{il}_1 \cos\theta z_0 \\ \mathrm{G}(\mathring{\mathbf{r}} - \mathring{\mathbf{r}}_0) & = \mathrm{G}(\mathring{\mathbf{r}}) \exp \end{array}$$

since:

$$(eqn \ 2.29)$$

$$\int_{-\pi}^{\pi} \frac{d\phi}{2\pi} cosn\phi exp exp exp exp = i^{-n} Jn(l_2 sin\theta)$$

$$\int_{-1}^{1} \frac{dz_{o}}{\pi \sqrt{1-z_{o}^{2}}} cospve^{-il_{1}cosvz} c = \int_{0}^{\pi} \frac{dv}{\pi} cospve^{-il_{1}cos\theta cosv}$$
$$= i^{-p}Jp(1_{1}cos\theta)$$

and with z=cosv and the facts that  $k_{\phi} \rightarrow (1-Z^2)^{4/2}, K_z \rightarrow (1-z^2)^{4/2}$  on the edges of the cylinder

$$K_{zn}(z) = \frac{1}{\pi} \int_{0.2}^{\infty} K_{z,n}^{p} \sin(p+1)v$$
 (eqn 2.31)

$$K_{\phi n}(z) = \frac{1}{\pi \sin v} \sum_{p=0}^{\infty} K_{\phi,n}^{p} \cos pv$$
 (eqn 2.32)

where  $K^p_{Z\,,\,n}$  and  $K^p_{\varphi\,,\,n}$  the sum of the inside and outside currents from equation 2.16

$$K_{z,n}^{(+)p} = K_{z,n} + K_{z,-n}$$
 (eqn 2.33)

$$K_{\phi,n}^{(+)p} = K_{\phi,n} + K_{\phi,-n}$$
 (eqn 2.34)

The final equations for the general case are:

$$E_{n}(r,\theta,\phi) = 0 (eqn 2.35)$$

$$-\frac{2i}{l_1 l_2 \xi_0 G(\overset{+}{r})} E_{\theta}(r,\theta,\phi) = -\sum_{n=0}^{\infty} \sum_{p=0}^{\infty} i^{-(n+p)} \frac{(p+1) \sin \theta}{l_1 \cos \theta}$$
(eqn 2.36)

$$\begin{split} & J_{p+1}(1_1 cos\theta) J_n(1_2 sin\theta) [K_{2n}^{(+)} P cosn\phi + iK_{2n}^{(-)} P sinn\phi] \\ & + \sum\limits_{n=1}^{\infty} \sum\limits_{p=0}^{\infty} i^{-(n+p)} \frac{n}{1_2 sin\theta} J_p(1_1 cos\theta) \mathbb{J}_n(1_2 sin\theta) [K_{\not p'n}^{(-)} P cosn\phi \\ & + iK_{\phi n}^{(+)} P sinn\phi]. \end{split}$$

$$-\frac{2i}{l_1 l_2 \xi G(r)} E_{\phi}(r,\theta,\phi) =$$

(eqn 2.37)

$$\sum_{n=0}^{\infty} \sum_{p=0}^{\infty} i^{-(n+p)} J_p(1_{2}\cos\theta) J_n^{1}(1_{2}\sin\theta) [-K_{\varphi n}^{(-)p} \sin n\varphi + iK_{\varphi n}^{(+)p} \cos n\varphi].$$

Since  $E_{p}(r,\theta,\phi)=0$  in the far field

(eqn 2.38)

$$E_{x}(r,\theta,\phi)=E_{\theta}(r,\theta,\phi)\cos\theta\cos\phi-E_{\theta}(r,\theta,\phi)\sin\phi$$

(eqn 2.39)

$$E_{y}(r,\theta,\phi) = E_{\theta}(r,\theta,\phi)\cos\theta\sin\phi + E_{\phi}(r,\theta,\phi)\cos\phi$$

(eqn 2.40)

$$E_z(r,\theta,\phi) = -E_{\theta}(r,\theta,\phi)\sin\theta$$

In the special case, with the cylinder positioned head-on to the antennas, the incident and scattered fields can be described as follows:

 $\theta = \pi$  (along the Z axis)

The scattered fields are:

(eqn 2.41)

$$\frac{-2\text{i}}{\text{l}_1\text{l}_2\xi_\circ\text{G}(\overset{+}{r})}\text{E}_\theta\left(\text{r},\pi,\phi\right) = \frac{1}{2}\sum_{p=0}^{\infty} \text{i}^{-(p+1)}\text{Jp}(\text{l}_1)[\text{K}_{\phi1}^{(-)p}\text{cos}\phi+\text{iK}_{\phi1}^{(+)p}\text{sin}\phi]$$

(eqn 2.42)

$$\frac{-2\mathrm{i}}{\mathbf{1}_{1}\mathbf{1}_{2}\xi_{\circ}G(\overset{\star}{r})}E_{\varphi}(\textbf{r},\pi,\phi)=\frac{1}{2}\sum_{p=0}^{\infty}~\mathbf{i}^{-(p+1)}Jp(\mathbf{1}_{1})[-K_{\varphi}^{(-)p}sin\phi+\mathbf{i}K_{\varphi}^{(+)p}cos\phi].$$

or in rectangular coordinates, along the -z axis:

(eqn 2.43)

$$\frac{-2 \text{i}}{ {1 \choose 1} {1 \choose 2} \xi_0 G_T^{'}} E_X(z) = \frac{1}{2} \sum_{p=0}^{\infty} \text{i}^{(p-1)} J_p(1_1) K_{\phi 1}^{(-)p}$$

(eqn 2.44)

$$\frac{-2i}{1_{1}1_{2}\xi_{0}G(\mathring{r})}E_{y}(z) = \frac{1}{2}\sum_{p=0}^{\infty} i^{-p}J_{p}(1_{1})K_{\phi 1}^{(+)p}$$

$$E_{z}(z) = 0.$$

The configuration of the cylinder and the incident field is shown in Figure 2.4

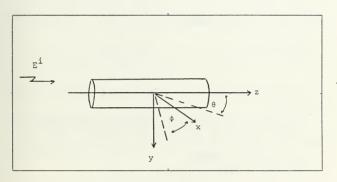


Figure 2.4 The Cylinder and the Incident Field

With the axis of the cylinder chosen along the z-direction and the incident field along the z axis, the incident field is given as:

$$\dot{E}^{inc} = \hat{e}_2 l^{ikz}$$
 ,  $\hat{e}_2 = -\hat{x}$  (eqn 2.46)

$$E_{\phi}^{inc} = +\sin\phi e^{ikz}$$
 (eqn 2.47)

Because  $E_{\varphi}^{\mbox{inc}}$  is an odd function in  $\varphi$ 

$$K_{\phi n}^{(+)}(z) = 0$$

From equations 2.43- 2.45 Ey(Z)=0,and  $\vec{E}$  (Z)= $\hat{x}F_X(Z)$  along the -Z axis.

Cross-section and phase of a scatterer are defined with a linearly polarized plane incident wave on the scatterer. The incident wave has unit strength and zero phase at the center of the scatterer.

The cross-section is given by:

$$\sigma = \lim_{r \to \infty} 4\pi r^2 \left| E^{SC} \right|^2$$
 (eqn 2.48)

and the phase shift is given by:

$$\delta = \arg(l^{-ikr}E^{SC}) = \arg[\sum_{p=0}^{\infty} i^{p}J_{p}(l_{1})K_{\phi 1}^{(-)p}]. \qquad (eqn 2.49)$$

since:

$$|E^{S_{q}^{C}}| = |\frac{1}{1}\frac{1}{2}\frac{1}{2}\frac{\xi_{0}G(r)}{\mu_{1}} \sum_{p=0}^{\infty} i^{p-1}J_{p}(1_{1})K_{\phi 1}^{(-)p}|$$
 (eqn 2.50)

the cross-section of the finite cylinder is:

$$\sigma = \lim_{\substack{\eta = 1 \\ r \to \infty}} \frac{\eta_{\pi r}^{2}}{\eta_{\pi r}^{2}} \left| \frac{1}{\eta_{\pi r}^{2}} \frac{1}{\eta_{\pi r}^{2}} \sum_{p=0}^{\infty} i^{p-1} J_{p}(1_{1}) K_{\phi 1}^{(-)p} \right|^{2}$$

$$= \pi a^{2} \left| \begin{array}{cc} \frac{kh\xi}{a} & \sum\limits_{p=0}^{\infty} i^{p-1} J_{p}(1) K_{\phi 1}^{(-)p} \right|^{2}$$

#### III. MEASUREMENTS AND RESULTS

Radar cross-section estimation is as much art as science. The air of mystery that surrounds it will only be removed by an increase in understanding of how a target scatters energy incident upon it.

A complete knowledge of the scattering behavior is available only for few bodies. For these bodies we do not have exact solutions for their radar cross section, the best we can do is to provide approximate values. Such values should be checked against experimental measurements.

This chapter describes some CW step frequency cross-section measurements carried out at the Naval Postgraduate School, including the experimental setup, the measurement procedure and the measurements results.

Computer programs for calibration of the experimental setup and target measurements process are given in Appendix A. An explanation to the programs can be found in previous thesis from the Naval Postgraduate School [Ref. 15].

#### A. LABORATORY DESCRIPTION

## 1. System Configuration

 $\label{the continuous} The physical setup of the \ laboratory can be divided into five parts:$ 

- -The RF source.
- -Transmitting and receiving antennas.
- -The anechoic chamber.
- -Target mount.targets.
- -Control and data processing equipment.

This setup is called a Radar range geometry and is designed for cross section measurements of models. The configuration of the entire system is shown in Figure 3.1.

A variety of radar range geometries have been developed during the years. Those radar range geometries were characterized by the way they have been designed to eliminate unwanted signals reflected from the foreground and background [Ref. 10]. An important component of the setup at the NavalPostgraduate School is the anechoic chamber [Ref. 11]. The anechoic chamber is used to approximate free space conditions in a closed environment. The anechoic chamber is enclosed with aluminium plates and internally lined with a radio frequency absorbing material. absorbing material provides the necessary attenuation to the reflections from the walls, floor and ceiling, and the aluminium surface provides protection against external sources of noise such as atmospheric noise, man made noise (radio, television, radar etc) and weather conditions. The characteristics of the absorber material are specified in Appendix B.

The target is supported by a styrofoam stand. The reason for choosing the material, is to minimize coupling between the stand and the target. [Ref. 12].

The radar cross section range utilizes two identical horn antennas to approximate a back scattering system. Both antennas are horizontally polarized. The antennas are mounted on a removable panel located in the front wall of the anechoic chamber. The antennas can be adjusted in all three axis, providing beam steering towards the target. The basic antennas characteristics are given in Table I, and the configuration of the anechoic chamber is shown in Figure 3.2

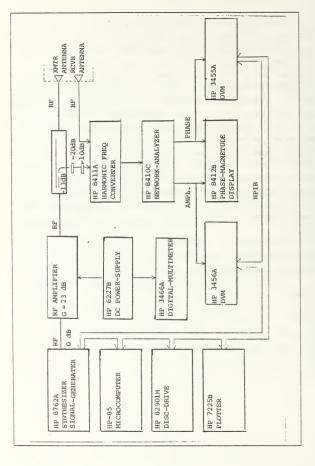


Figure 3.1 Block Diagram of the System

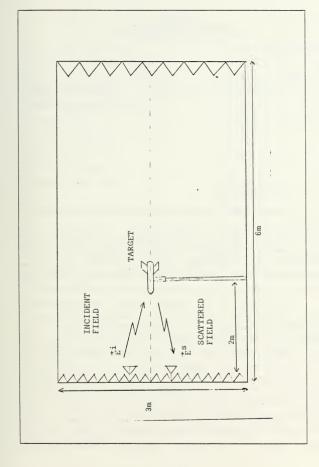


Figure 3.2 Configuration of the Anechoic Chamber

# TABLE I

Frequency range	4-18 GHz		
Gain	6-12dB		
VSWR (max)	3.2:1		
Isolation between cross polarization	<20dB below 5.5 GHz		

>20dB above 5.5 GHz

## 2. Instrumentation

 $\label{the equipment of the experimental setup is listed in $$ Table II .$ 

The RF signal to the system is provided by the signal generator (HP-8672A). The RF output from the signal generator is amplified by an RF amplifier (Avantek-SA-83-2953). (The DC-power supply,HP-6227B,to the amplifier is monitored by a digital multimeter,HP-3466A).

The amplified RF signal passes the directional coupler (Narda-5292) and goes to the transmitting antenna. The directional coupler has a coupling coefficient of -13 dB and provides the reference signal to the harmonic frequency converter (HP-841LA) after the signal is further attenuated by 30 dB. The scattered signal from the target goes through the receiving antenna to the first port of the harmonic frequency converter. The frequency converter and the network analyzer (HP-8410C) convert the test channel signal into two 278 KHz signals containing the magnitude and phase information of the test channel signal relative to the refence channel signal. Both signals enter the phase-magnitude

TABLE II
LIST OF EXPERIMENTAL DEVICES

TTD 0 F

-Microcomputer	HP-85
-Synthesized Signal Generator	HP-8672A
-Harmonic Frequency converter	HP-8411A
-Network Analizer	HP-8410C
-Phase-Magnitude Display	HP-8412B
-RF Amplifier	Avantek sa-83-2953
-DC Power supply	HP-6277B
-Digital Multimeter	HP-3466A
-Digital voltmeter	HP-3455A
-Digital voltmeter	HP-3456A
-Flexible Disc Drive	HP-82901M
-Plotter	HP-7225B
-Directional Coupler	Narda 5292

display (HP-8412B) The DC plotter outputs of the display unit are fed to the two DVM's (HP-3455A,HP-3456a).

The complete setup is controlled and the data is processed by the microcomputer (HP-85) and the results are stored on discs.

## 3. Targets

The most direct means of obtaining knowledges about radar cross sections are by measurement of the radar return from the target itself or from an accurate model of the target [Ref. 13]. One advantage of a radar range is the practicability of testing models that are smaller and cheaper than full scale targets. The radar wavelength is scaled by the same factor as are the dimensions of the model. If D is any given dimension of the target, and D is the equivalent dimension of the model, the following scaling relation is employed:

$$D_{M}/D_{O} = \lambda_{M}/\lambda_{O}$$
 (eqn 3.1)

where  $\lambda_{M}$  is the wavelength used for the measurement and  $\lambda_{0}$  is the wavelength for which the target radar cross-section is required.

At the same time, the measured cross section is altered in proportion to the change in power captured as a result of dimensional changes. If  $\sigma_M$  is the observed cross section of the model, the target cross section  $\sigma_0$  is given by:

$$\sigma_0 = \frac{\lambda_0^2}{\lambda_M^2} \sigma_M \qquad (eqn 3.2)$$

Exact scaling requires the model conductivity to be equal to the target conductivity multiplied by the ratio  $(\lambda_0^{}/\lambda_M^{})$ , and model permittivity and permeability at the test frequency to equal corresponding target electrical properties at the operational frequency. It is necessary to make sure that the surface conductance of the model won't be smaller than the target.

The targets which are tested in the measurements are thin walled tubular cylinders made of brass. There are 20 cylinders of various lengths and diameters. The dimensions of the targets are given in Table III .

For the calibration of the system a 3.187 inch aluminium sphere is used.

Some measurements are taken with a cylinder with fins attached . The description of the cylinder with fins is shown in Figure 3.3 .

TABLE III
TARGET-CHARACTERISTICS

TARGET		LENGTH	DIAMETER	THICKNESS
cylinder	1	2.00"	0.375"	0.012"
cylinder	2	2.00"	0.500"	0.014"
cylinder	3	2.00"	0.750"	0.011"
cylinder	4	2.25"	0.375"	0.012"
cylinder	5	2.25"	0.500"	0.014"
cylinder	6	2.25"	0.750"	0.011"
cylinder	7	2.50"	0.375"	0.012"
cylinder	8	2.50"	0.500"	0.014"
cylinder	9	2.50"	0.750"	0.011"
cylinder	10	2.75"	0.375"	0.012"
cylinder	11	2.75"	0.500"	0.014"
cylinder	12	2.75"	0.750"	0.011"
cylinder	13	3.00"	0.375"	0.012"
cylinder	14	3.00"	0.500"	0.014"
cylinder	15	3.00"	0.750"	0.011"
cylinder	16	1.50"	0.375"	0.012"
cylinder	17	4.50"	0.750"	0.011"
cylinder	18	2.50"	0.625"	0.014"
cylinder	19	3.75"	0.750"	0.011"
cylinder	with fins	3.00"	0.750"	0.011"

Figure 3.3 Tubular Cylinder with Fins

#### B. MEASUREMENT PROCEDURE

## 1. Calibration of the System

 $\hspace{1.5cm} \text{As a first step in the measurement, a calibration of } \\ \text{the system must be done.}$ 

To calibrate the system output, a target of known cross section (usually a metal sphere) is placed at at the target support to fix the level of the calibration curve. This measurement assures that the entire system is calibrated in the proper frequency range.

Measurements are taken at many fixed frequencies between 10 to 15 GHz. Since both the transmitting and receiving antennas are horizontally polarized, the measurements are in one dimensional plane.

The calibration of the system is divided into two steps:

- Take measurements without target in the anechoic chamber. (background data).
- 2.Inserte a 3.187 inch diameter aluminium sphere and repeat the measurements.

With the target in place, the received signal is a vectorial sum of the target echo and the background radiation. By taking background data, direct back scattering from the targets support and the walls of the anechoic chamber can be substracted from the vector sum.

The quality of the calibration of the system is tested and checked by comparing a new set of measured data on the sphere to thier theoretical values whenever a calibration process is done.

After the calibration of the system is accomplished, measurements of target cross sections can be started.

#### 2. Measurements of the Targets

The measurements of the cylinders are taken in the frequency range of 10 to 15 GHz in steps of 0.1 GHz. The decision to limit the frequency range inspite of the ability of the equipment to operate beyond this range is due to the following reason:

Only in this frequency range consistent data can be obtained through averaging the data obtained from several frequency scans.

All the measurements are taken while the cylinders are positioned "Head-on" to the antennas as seen in Figure 3.4 A description of the cylinders tested is given in Table III, and reasons for choosing those particular dimensions for the cylinders are given in Chapter .

As a rule of thumb, far field approximations are good under the following conditions:

$$r > 10\lambda$$
 (eqn 3.3)

$$r > \frac{2D^2}{\lambda}$$
 (eqn 3.5)

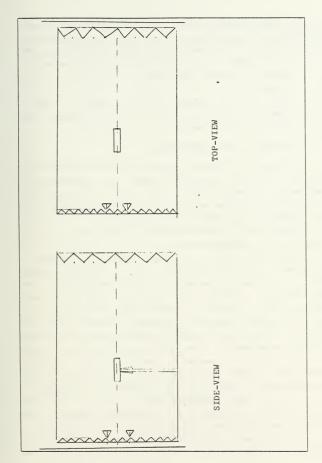


Figure 3.4 The Geometry Range

where r is the distance between the target and the antennas (receiving and transmitting), D is the largest dimension of either the target or the antennas and their separations, and  $\lambda$  is the wavelength. In our case, all conditions are met.

The experimental results for the targets in term of plots of their cross section and phases versus frequency are given at the end of the chapter.

## 3. Sources of Measurement Errors

Recognition of error sources in the measured data is difficult. Only in few special cases, it is possible to recognize the presence of an error, and to determine its source by observing the deviation of a target cross section versus frequency plot from anticipated behavior.

The errors in the measurement in our case are due to system noise and background noise.

The system noise, is caused mainly by the receiver. The network analyzer (HP-8410C) is a harmonic mixing receiver. It selects a harmonic of its internal VCO for the local oscillator frequency used to down convert the test frequency to the first IF. Harmonic skip errors can occur when the receiver selects a different harmonic (and VCO frequency) for the same frequency between system calibration and target measurement. The local oscillator power varies from one harmonic to another, thus the mixer transfer-characteristic varies. This fact causes random magnitude and phase variations of from 0.1 to 0.4 dB and up to 2 degrees.

As for the background noise, it is caused by coupling between the target and its mechanical support, and by strong coupling between the antennas in low frequencies.

In our case, the coupling between the antennas is reduced by going to higher range of frequencies. The

coupling between the target and its mechanical support is unavoidable and finally the receiver noise is reduced by averaging the measured results.

Figure 3.5 Cylinder 1: Cross-section

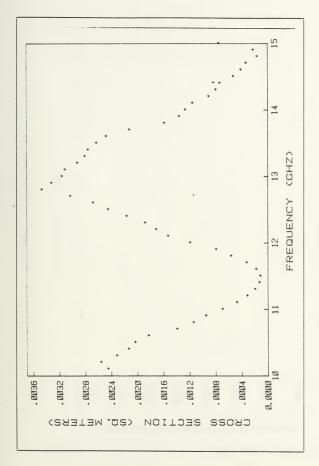


Figure 3.6 Cylinder 2: Cross-section

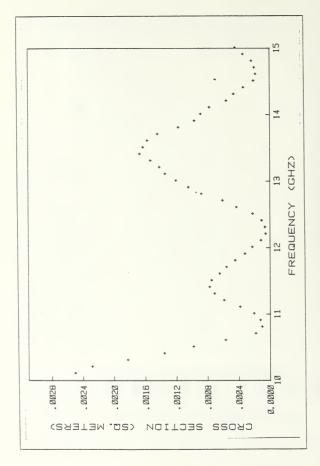


Figure 3.7 Cylinder 3: Cross-section

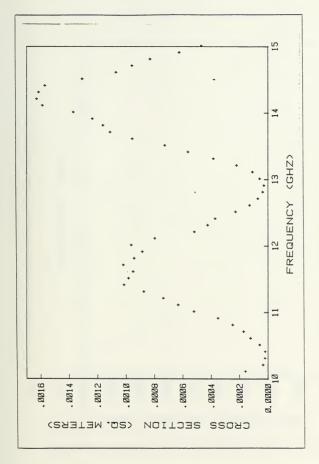


Figure 3.8 Cylinder 4: Cross-section

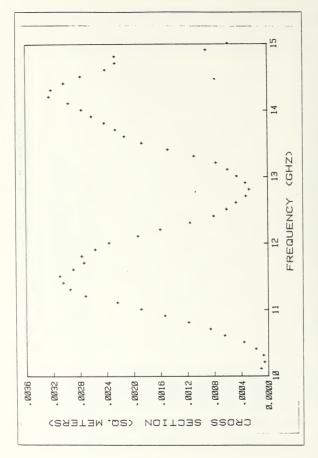


Figure 3.9 Cylinder 5: Cross-section

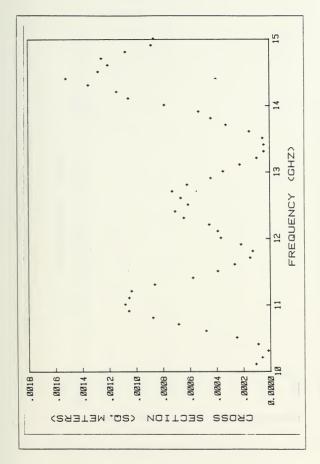


Figure 3.10 Cylinder 6: Cross-section

Figure 3.11 Cylinder 7: Cross-section

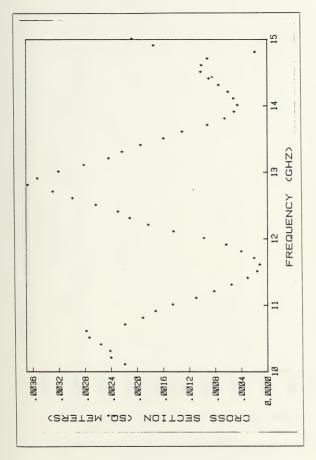


Figure 3.12 Cylinder 8 : Cross-section

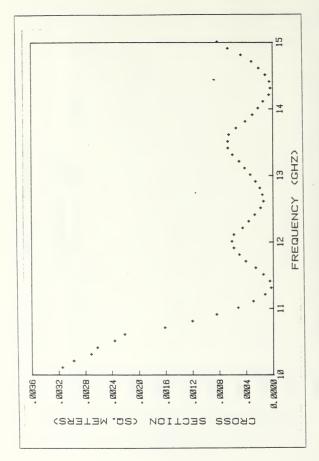


Figure 3.13 Cylinder 9: Cross-section

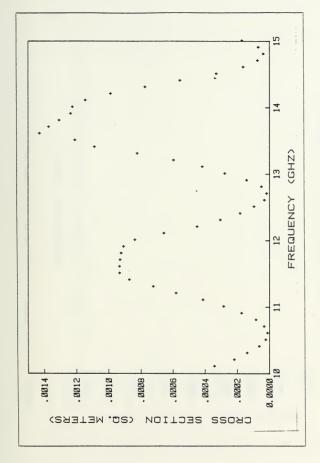


Figure 3.14 Cylinder 10 : Cross-section

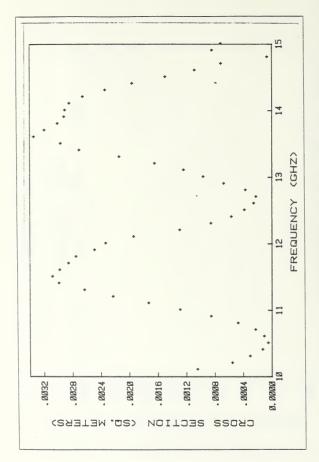


Figure 3.15 Cylinder 11 : Cross-section

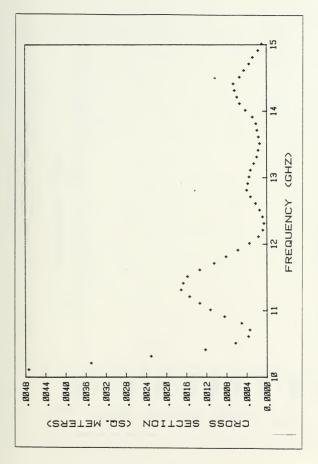


Figure 3.16 Cylinder 12: Cross-section

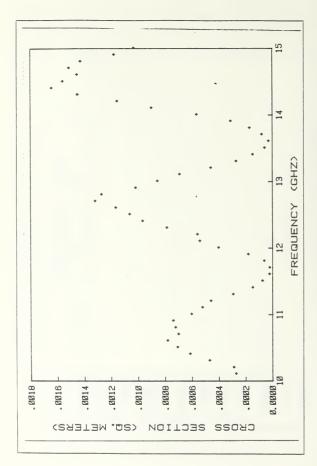


Figure 3.17 Cylinder 13 : Cross-section

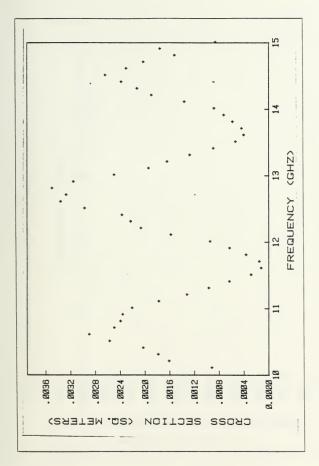


Figure 3.18 Cylinder 14 : Cross-section

Figure 3.19 Cylinder 15: Cross-section

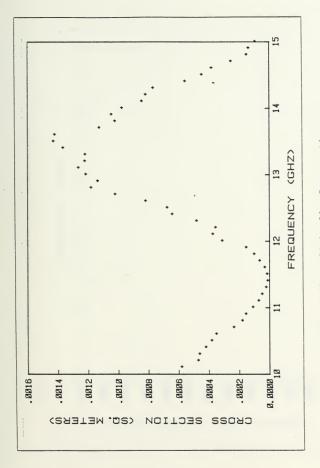


Figure 3.20 Cylinder 16: Cross-section

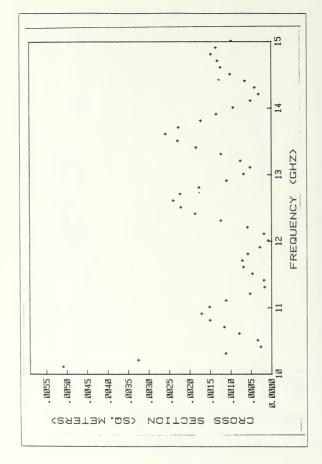


Figure 3.21 Cylinder 17: Cross-section

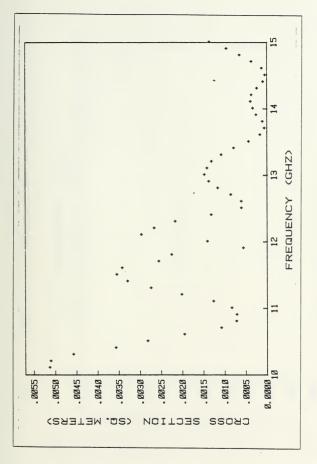


Figure 3.22 Cylinder 18 : Cross-section

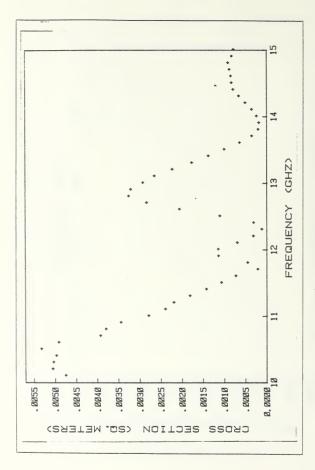


Figure 3.23 Cylinder 19 : Cross-section

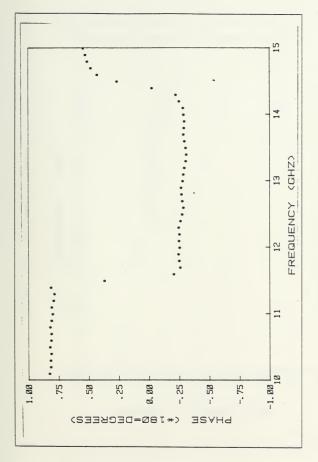


Figure 3.24 Cylinder 1 : Phase

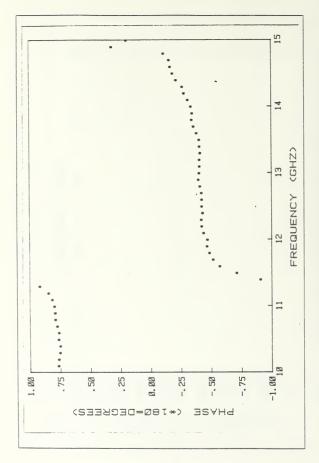


Figure 3.25 Cylinder 2: Phase

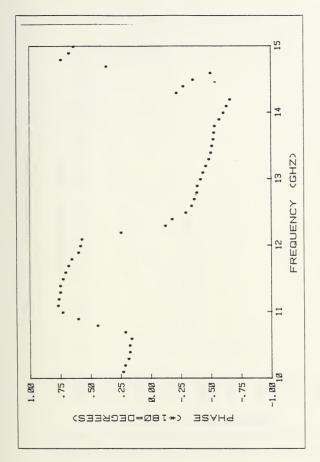


Figure 3.26 Cylinder 3: Phase

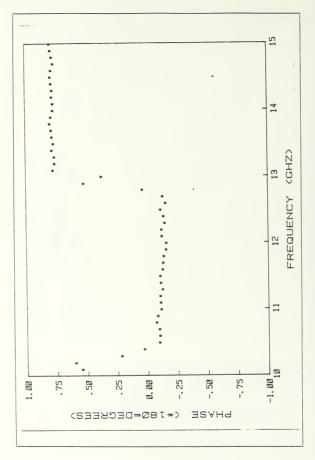


Figure 3.27 Cylinder 4: Phase

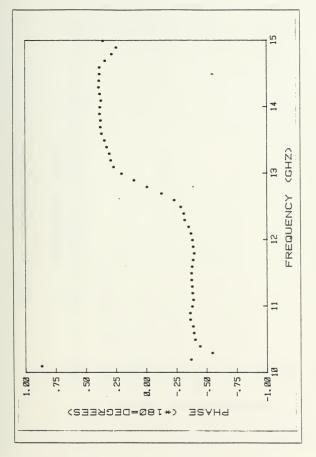


Figure 3.28 Cylinder 5: Phase

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Figure 3.29 Cylinder 6: Phase

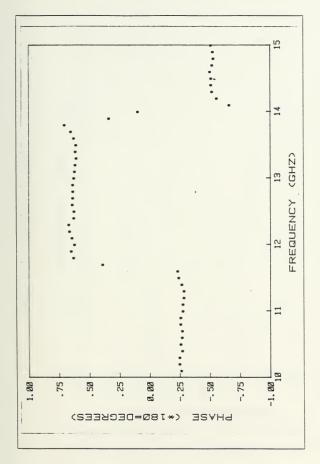


Figure 3.30 Cylinder 7: Phase

Figure 3.31 Cylinder 8 : Phase

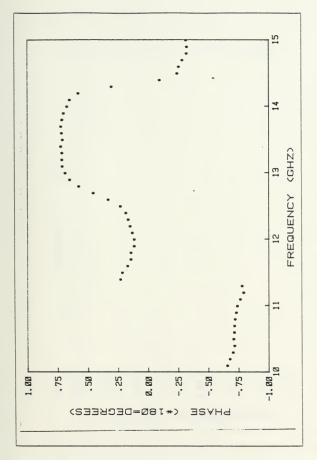


Figure 3.32 Cylinder 9 : Phase

Figure 3.33 Cylinder 10 : Phase

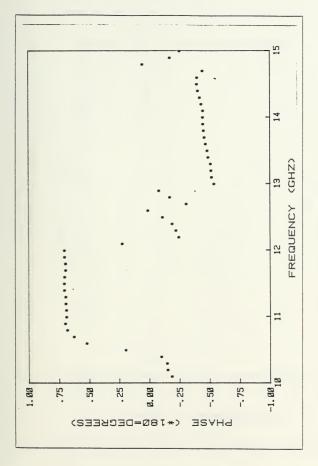


Figure 3.34 Cylinder 11 : Phase

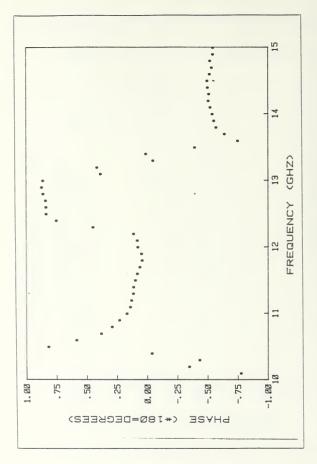


Figure 3.35 Cylinder 12: Phase

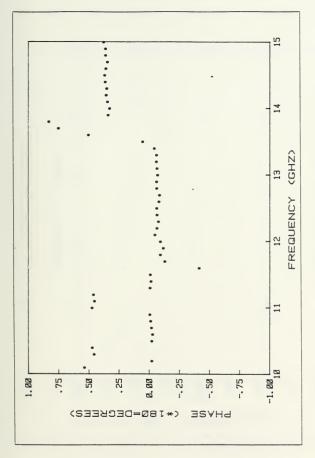


Figure 3.36 Cylinder 13 : Phase

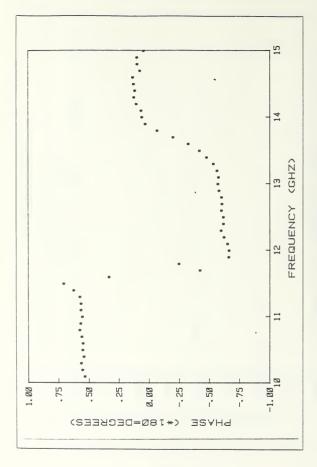


Figure 3.37 Cylinder 14: Phase

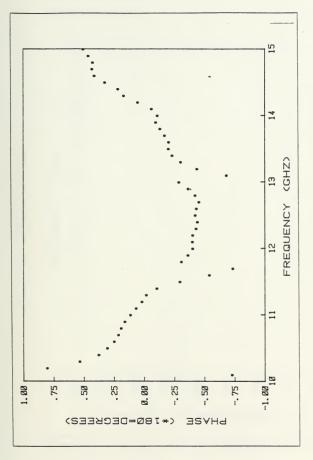


Figure 3.38 Cylinder 15: Phase

Figure 3.39 Cylinder 16: Phase

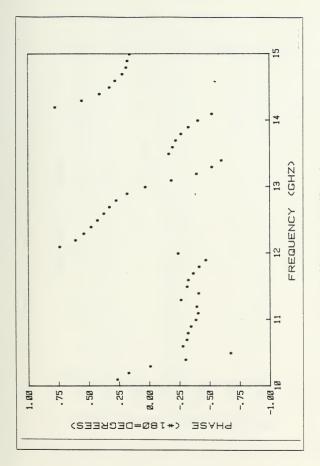


Figure 3.40 Cylinder 17: Phase

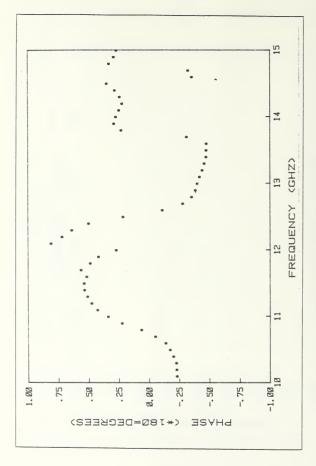


Figure 3.41 Cylinder 18: Phase

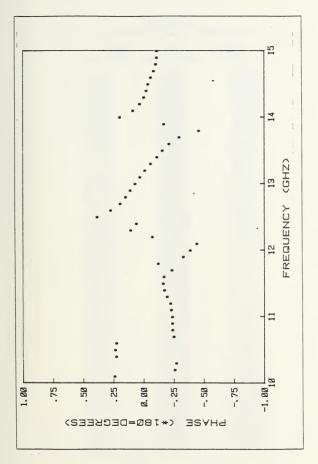


Figure 3.42 Cylinder 19 : Phase

TABLE IV

## CYLINDER 1 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG /180)
10000000000000000000000000000000000000	7976683470255941042657777979777779777668534770253465547776666534770255346554776666534770255346554776666547776666554770255346567666554770255346567666554770255346567666564776665547702553466666666666666666666666666666666666	811273757578585828149910898769997777777777777777777777777777777

TABLE V
CYLINDER 2 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)		(DEG./180)
10000000000000000000000000000000000000	######################################	7.4469403391132178339987746626351498146777781144832488484848485448544854485485391831777314483248848484848484848484848484848485384884848484

TABLE VI
CYLINDER 3 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
10000000000000000000000000000000000000	0921304 0921304 0921304 0921304 0931304 0931305 093130	1851757497851795578887858878579497487878788888888878688785788878888888888

TABLE VII
CYLINDER 4 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQP.METER)	PHASE (DEG./180)
1906-1906-1906-1906-1906-1906-1906-1906-	00015 000001 000000	5389263993859938570153149562523518837725529

# TABLE VIII CYLINDER 5 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHŽ)	(SaŘ.MĚTER)	(DEG./180)
100 000 000 000 000 000 000 000 000 000	.000003 .000005 .00000005 .0000005 .0000005 .0000005 .0000005 .0000005 .0000005 .00000000	- 14949597479991897294897919497959597949794979497949794979497949794

TABLE IX

CYLINDER 6 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SOR.METER)	(DEG./180.
19999999999999999999999999999999999999	99999999999999999999999999999999999999	230144 351926 -73113244 -77132444 -771344 -771344 -771344 -771344 -771344

TABLE X
CYLINDER 7 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
10000		10100
(GHZ)	(SQR.METER)	(DEG /180)
	0.0054	
10 10	00051	27541
10.20	.00062	· - 26286
10.20	.00002	. 20200
10.30	.00063	*26206 26263
10.49	99969	28534
		. 20007
10.50	. 00000	27050
19.69	99985	28241 28484
		. 20271
10.70	.00069	28484
10.80	.00062	26957 27013
		. 20001
10.90	.00058	27013
11.00	.00054	- 28944
14.40		00007
11.10	.00041	- 28903
11.20	.00031	- 28944 - 28903 - 29767
11.30	.00023	20071
11.50	. 6665	29871
11.40	.00013	27982
11.50	.00004	29871 27982 25442
11.06		- 26263 - 28534 - 27050 - 28484 - 26957 - 289484 - 26957 - 289483 - 28963 - 29767 - 27982 - 25442 - 24695
11.60	. 99991	24695 .37470
11.70	. 00000	37470
11.12		. 51 415
11.80	. 00004	.61601
11 90	.00009	63519
10.00	00020	60815
12.00	00020	.60815
12.10	.00036	. 62896
12 20	.00055	.64955
16.55		. 9 = 2 0 0
12.30	.00060	. 65435
12 40	.00070	.61442 .61292
40.50		54000
12.00	.00090	.61252
12 60	.00105	. 62563
10.70	00118	63120
12.70	. 66110	.62563 .62120 .62072
12.89	.00137	.62072
12 90	.00137	.62265
12.50		.64460
13.00	.00123	.61241
17 10	.00109	.60690
12.15	. 50165	. 66620
13.20	.00096	.60066
17 70	.00086	.59325
10 00		. 95929
13.40	.00070	.59960
17 50	.00055	. 59522
12.20		.03022
13.60	. 88848	.61393
17 70	.00023	.63828
12.15		.00020
13.80	.00009	.69122
111000 0000000000000000000000000000000	99993	.69122 .32204
14.00		00707
14.00	.00001	.08397 67478
14.16	.00005	67478
14.20	00017	- 57053
14.30	.00031	53141
14.40	00049	52406
		22400
14.50	.00070	- 52495
14.60	00087	- 51447
	.00102	53291
		53291
14 80	.00128	- 54296
14.90	.00149	
		7.20020
15.00	.00160	53833 52256

TABLE XI

#### CYLINDER 8 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SOR.METER)	(DEG./180)
00000000000000000000000000000000000000	789931669993792839887495148971666587948927892888622822211815288398862388625887918685887918868444977278989662886868686868686868686868686868686	- 1000809517.2449700171100054105420555555555555555555555555555

TABLE XII
CYLINDER 9 : MEASUREMENT DATA

FPEQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
09000000000000000000000000000000000000	032233999999999999999999999999999999999	- 67984427-987915112299999999999999999999999999999999

TABLE XIII

### CYLINDER 10 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
193696666666666666666666666666666666666	99931 99931	- 98947-981-98-99-99-99-99-99-99-99-99-99-99-99-99-

#### TABLE XIV

# CYLINDER 11 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQP.METER)	(DEG./180)
100 000 000 000 000 000 000 000 000 000	238178811576611117666757712475346835554222284842799.8225144694429 0000000000000000000000000000000000	2013391114432179722116337752021163274465977511866217474744597752175311862174474646961746747446597797217557722991018622774746597792175577722755116497474445977792175577224575116497474445977921755772245751164974744459779217557722457511649747444597797217577722991018672175772245751722457474445474445977972175772457511649747444514144541445977972175772457511649747444541445977974744454144541444597797474445414459779747444541445977444454144597744445414459774444541445977444454144597744445414459747444541445974744454144459747444541444444

TABLE XV
CYLINDER 12 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
10 10 10 10 10 10 10 10 10 10 10 10 10 1	0474 0474 0474 04727 091199 046737 046737 046737 046737 046737 04674 04757 047	- 799866 - 4443157 - 4443157 - 4443157 - 9502942 - 9502942 - 9502942 - 1524365 - 152436 - 1524365 - 152436

TABLE XVI
CYLINDER 13 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
19000000000000000000000000000000000000	002366000000000000000000000000000000000	115565561988997388991446457374958894246988746948444497699674756987469484444976996747569674769694769487694876948769696747569876769876769877657756775677566876769876769876765776698776766987767669877677669877677669877677669877677669877677669877677669877677669877677669877677669877677669877677669877677677677677767

TABLE XVII
CYLINDER 14 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
00000000000000000000000000000000000000	9986001700447773885944166991533004477114461813770587777388594116699153300447777388594116699153300447711446177138859415629044733885941709604960496049533009604979738994477389944773899499999999999999	- 177349 - 1773

TABLE XVIII
CYLINDER 15 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
10 120 10 10 10 10 10 10 10 10 10 10 10 10 10	9840 98965 98149 98278 98331 98438 98331 98438 98249 98123 98438 98123 98438 98141 98438 98438 98444	- 742544 7519244 751926967 7519269687 7519269687 7519269687 75192697 75192697 75192697 75192697 751927 7519

TABLE XIX
CYLINDER 16 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
1000 400 400 400 400 400 400 400 400 400	090445 090445 090445 090445 090447 0908447 090806 09080	- 1132368846716935996273513269716935969716935969716935976973597697359769737359769737359769737359769737766574887776657488777663714847746371847766371484777663714848977766371884897766371884897776637188489777663718848977766371884897776637188489777766371884897776637188489777663718878897778897889789789789789789789789789

TABLE XX

CYLINDER 17 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG /180)
10	9594 9854 98198 98198 98198 98197 981146 981146 981447 98147 98147 98147 98147 98147 98147 98147 98147 98147 98147 98147 9	25118 25118 1621838 1621838 1621839 1621839 1631839

TABLE XXI
CYLINDER 18 : MEASUREMENT DATA

FREQ (GHZ)	CRSEC (SQR.METER)	PHASE (DEG./180)
11111111111111111111111111111111111111	781993878599719998876299178192369914481599938784993887629913966716488876299139699999999999999999999999999999999	22197477 2219747 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 221974777 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 2219747 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 22197477 221974777 221974

TABLE XXII

#### CYLINDER 19 : MEASUREMENT DATA

FREQ	CRSEC	PHASE
(GHZ)	(SQR.METER)	(DEG./180)
988 998 988 988 988 988 988 988 988 988	00338 .00231 .00221 .00136 .00136 .00292 .00292 .00292 .00292 .00297 .00056 .00057 .00057 .00031 .00112 .00137 .00137 .00137 .00124 .00124 .00127	- 2439666767220110000100010001000100010001000100010

#### IV. SUMMARY AND CONCLUSIONS

#### A. ANALYSIS OF MEASURED DATA

The back scattering cross section of a tubular cylinder depends upon several parameters. To simplify the problem, many of them were kept constant: The polarization of the receiving and transmitting antennas were kept constant during all measurements, all the cylinders were from the same material and the aspect angle and tilt of the different cylinders relative to the antennas plane were unchanged.

The parameters that were changed and their effects on the cross section were the subjects of this study were:

- 1. The cylinder length. (2h)
- 2. The cylinder diameter. (2a)
- 3. The transmitting frequency.

The cross section of a tubular cylinder can be written as a function of three parameters as shown in equsion 4.1

$$\sigma = F_1(a,h,f) = F_2(a,h,k)$$
 (eqn 4.1)

where:

$$k = 2\pi/\lambda = 2\pi f/c$$

$$c = 3 \times 10^8 \text{m/sec}$$

From the theory described in Chapter the following relations can be obtained:

(egn 4.2)

$$\frac{\sigma}{\pi a^2} = F_3(l_1, l_2) = F_4(l_1, \frac{l_2}{l_1}) = F_s(ka, \frac{h}{a})$$

where:

$$l_2 = kh$$

The measurements taken can be devided into two steps. In the first step, 19 cylinders with different lengths and diameters were placed in the anechoic chamber, and plots of cross section and phase versus frequency were obtained on each of the cylinders. (figures 3.5 through 3.42) Except for gathering data and viewing the resonance frequencies, there is nothing much to say on the results because too many parameters were changed and the resulting analysis became too difficult and complicated.

The second step of measurements was based on the conclusion from the theory (equation 4.2).

Only eight scaled cylinders were put into the anechoic chamber, four of the cylinders had a constant ratio of length to diameter equals 4, (h/a=4), and four cylinders had a constant ratio of length to diameter equals 6, (h/a=6).

An overall plot with expanded frequency range was achieved by using those scaled cylinders.

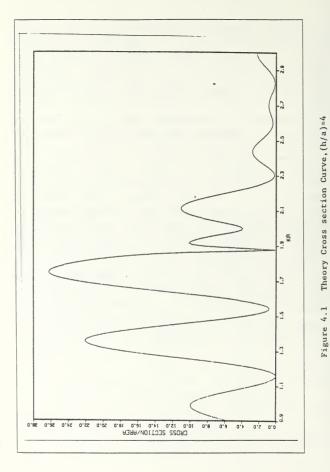
The overall plots can be compared to the theoretical plots obtained by Prof.H.M.Lee at the Naval Postgraduate School.

Both the theoretical and experimental plots of cross section/area and phase versus ka are shown in Figures 4.1 through 4.8 and the overlapping ranges of the scaled cylinders in ka can be noticed from Tables XXIII through XXX The theoretical and experimental curves can be seen on the same graph for the two sets of scaled cylinders in Figures 4.9 and 4.10

Comparison of the experimental results to the theoretical plots show agreement away from the cutoff frequencies of the  $\rm H_{11}$  circular waveguide mode at ka=1.8415.

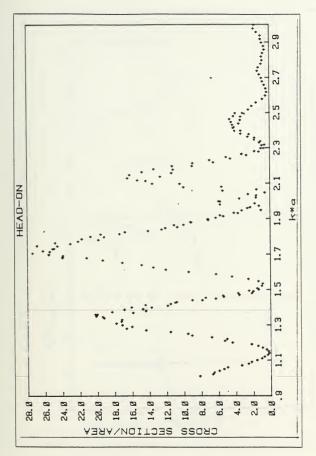
In the theoretical calculations an infinitesimal wall thickness is assumed for the tubular cylinder, but practically the cylinders used in the measurements had some thickness (Table III). The outer diameter of the cylinder is used for 2a in the computations while the  $\rm H_{11}$  mode cutoff frequency depends on the inner diameter of the cylinder. This fact caused the deviation in the cross section plots between the theoretical curve and the experimental curve near the cutoff frequency.

Problems occur in the phase plot because the averaging procedure did not properly take care of the phase shifts with values near  $\pm$  180 degrees. Because of the complicated behavior of the phase shift near the H  $_{11}$  cutoff at ka=1.8415 and near ka=2.4046, no conclusion can be made related to the actual behavior of the phase shift curve.

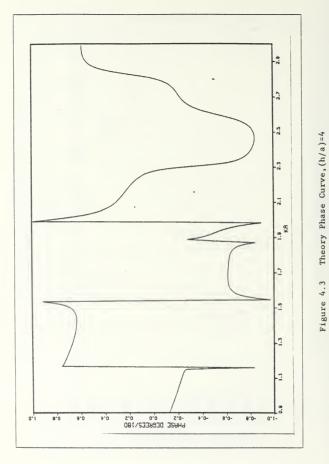


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Figure 4.1



Experimental Cross section Curve, (h/a)=4 Figure 4.2



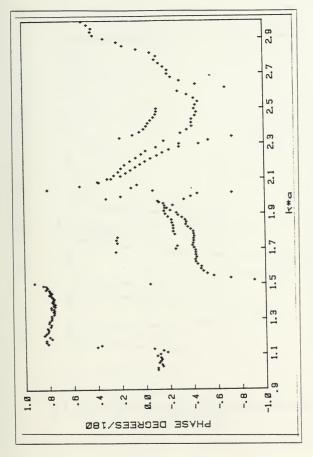


Figure 4.4 Experimental Phase Curve, (h/a)=4

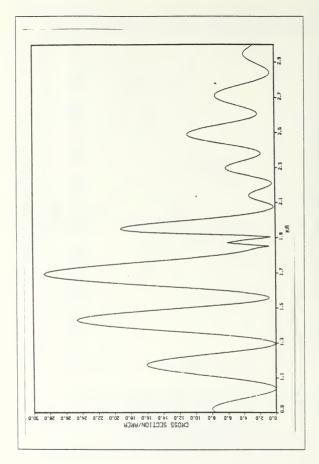
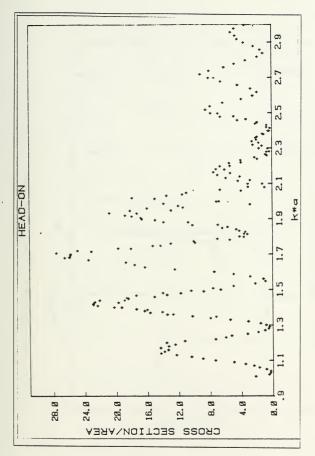


Figure 4.5 Theory Cross section Curve, (h/a)=6



Experimental Cross section Curve, (h/a)=6

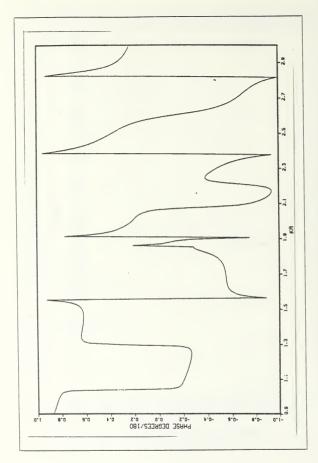


Figure 4.7 Theory Phase Curve, (h/a)=6

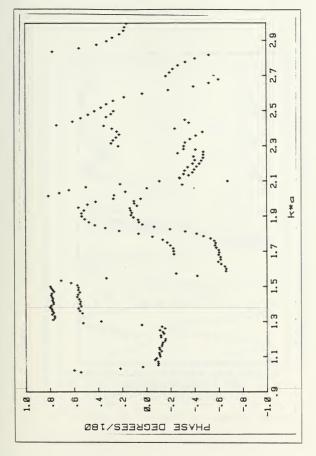


Figure 4.8 Experimental Phase Curve, (h/a)=6

Figure 4.9 Summary Cross section Curve, (h/a)=4

Figure 4.10 Summary Cross section Curve, (h/a)=6

## TABLE XXIII CYL16.4:1 RATIO

(K*A)	CRSZAREA	PHASE
1123456789981123456789981234567899812345678998123456789981234567899812345678998123456789981234567899812345678998123555555555555555555555555555555555555	981489 7647129 7647129 7647129 76471216	27921732826217321122298897887683148897766937148897769217321723792217323747766977759731112229889768876887769777697774888769877788877789787778977789

# TABLE XXIV CYL.2,4:1 RATIO

(K <b>#A</b> )	CRS/AREA	PHASE
11111111111111111111111111111111111111	9600174280 60200774280 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 196000077480 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 1960000774280 19600000774280 196000000000000000000000000000000000000	75482 746924 7466724 7466724 7466724 747682 7486924 774824 7778308 80011162 902167333 150219 902267333 150219 148484 148484 14848 14

TABLE XXV CYL.18,4:1 RATIO

(K#A)	CRS/AREA	PHASE
[.70 [.71 [.73 [.75 [.76 [.78 [.80	23 5 7 3 7 3 5 7 3 5 7 3 7 3 7 3 7 3 7 3	25796 -225797

TABLE XXVI CYL.15,4:1 RATIO

(K*A)	CRS/AREA	PHASE
9135791135791357913579135791357913579135	1. 39855 2. 244254 9. 74599 15. 244254 11. 12595 15. 241224 11. 12595 15. 241224 11. 12595 15. 241224 16. 25122 17. 26108 17.	- 748443888714858845878885888588585885885885885885885885885

## TABLE XXVII CYL.4,6:1 RATIO

(K#	(A) CRS	/AREA	PHASE
123745678981237456789812374567899812374567898 88888888888888881111111111111111111	11607 0377 1109 05 0377 05 4 00 6 7 4 00 4 00 4 7 9 6 7 7 7 10 6 6 6 6 7 7 7 7 10 6 6 6 6 7 7 7 7 7 7 7 7 7 7 7 7 7 7 7		53000 53

### TABLE XXVIII CYL.14,6:1 RATIO

TABLE XXIX
CYL.19,6:1 RATIO

(K <b>≭A</b> )	CRSZAREA	PHASE
117777688137568891755688917746891774689177468991774448991774468991774448991774468991774448991774468991774448991774468991774444444444444444444444444444444444	25 791589 221 894488 225 894488 227 894488 227 894488 228 894788 228 894788 229 894788 2317 895894 2418 895894 241	- 244388 - 243896 - 238966 - 238966 - 238966 - 238966 - 238966 - 238968 - 238968 - 238968 - 238910 - 419956 - 42988 - 48910 - 48928 - 236389 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 296897 - 376896 - 472926 - 488928 - 48928 - 327379 - 237389

TABLE XXX
CYL.17,6:1 RATIO

(K)	‡A) C	RS/AREA	PHA:	
(\$\) 1357911357911357913557913557913579 22222222232379344479 2222222222323793444791555559666677777888888869999	17. 11. 3.	RS 3000 257-2027 3000 3000 3000 3000 3000 3000 3000		8
2.97	4.5 3.2	7189 7827	. 1	6087 4278

#### B. CONCLUSIONS AND RECOMENDATIONS

The results showed agreement between the theory and the experimental data obtained. Some deviations between the theoretical and experimental plots could be explained.

The problem caused by averaging as observed in the phase plot can be overcomed by refining the program in the following way: If the two values to be averaged deviate more than a pre-determined value, say, 200 degrees, both will be converted to a positive value before the average is taken. The averaged value is then adjusted so that it lies between -180 and +180 degrees. A new set of measurements with targets having length to inner diameter ratios of 4 and 6 is planned. This refinment of averaging procedure will be adopted.

For the next step in the comparative study of target back scattering characteristics based on the canonical model of the tubular cylinder of finite length, more complicated models varying from a tubular cylinder to a missile by adding fins and wings will be constructed and studied. Figures 4.11 and 4.12 show the experimental patterns of a finite tubular cylinder with four fins in two aspect angles, the first, with the fins parallel to the antennas, and the second twisted in 45 degrees. Both Figures can be compared to Figure 3.19 showing the same cylinder without the fins.

This work is the first step to gain the capability of radar target identification. Its contribution lies in the justification that the canonical model is a useful one: the theory is correct so that the theoretical results can provide information about surface current distribution on the cylinder.

Another application of this work could be in the study of the back scattering characteristics of aircraft engines.

Since the signals are modulated, target detection will be easier. And because of the smaller size of the air-intake, the results will be directly appliable to radars now in existance.

Figure 4.11 Cylinder 15 with Fins, 90 Deg

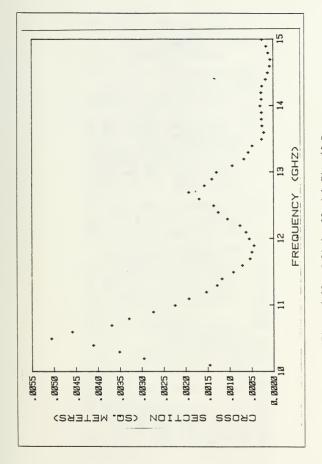


Figure 4.12 Cylinder 15 with Fins,45 Deg

### APPENDIX A COMPUTER PROGRAMS

```
10 | "SPHERE DRIVEO"
  20
 30
      COMPUTE BACK SCATTERRED
     1
 49 | FAR-FIELD FROM A FERFECTLY
      CONDUCTING SPHERE
 59
      THE INCIDENT FIELD IS A
 60
     ! LINEARLY POLARIZED PLANE W
 70
    AVE WITH ZERO PHASE AT THE ENTER OF THE SPHERE.
 80
    ! THE THEORETICAL VALUES TO
     BE COMPUTED ARE THE BACK-SC
    ATTERING CROSS-SECTION AND
    ! THE PHASE OF THE EAR FIELD
 90
      INTERPOLATED TO THE CENTER
    OF THE SPHERE.
100
      THEORETICAL VALUES ARE
110
      STORED IN FILES OF 800
120 1
    ! RECORDS, ONE FOR EACH FREQ
139
    UENCY FROM 2.02 GHZ TO 18 GH
      AT 0.02 GHZ STEPS.
149
159
      "THEORY DRIVEY" FOR THE
160
      1" DIAMETER SPHERE
179
      "THEOR3.ORIVE1" FOR THE
180
      3.187" DIAMETER SPHERE
190
      "THEOR4 DRIVE1" FOR THE
     4.75" DIAMETER SPHERE
200
210
      "THEORE DRIVE1" FOR THE
220
      6" DIAMETER SPHERE
230
240
    ! FILE H# STORES THE
    COMPUTED RESULT
260
    H≢="THEOR6.ORIVE1"
270
280
   A0=6*.0254/2 ! SPHERE
    RADIUS IN METERS
290
300
310 X9=2*PI ! PARAMETER
320
330
340 01=2 ! STARTING FRED IN GHZ
350 02=18 ! FINAL FREQ IN GHZ
360 04= 02 ! FREQ STEP IN GHZ
370 ON ERROR GOTO 390
380 PURGE HI
390 OFF ERROR
```

```
400 CREATE H$.S00.16 | OPEN A NE
     W FILE WITH 800 RECORDS
     ! OF 16 BYTES EACH, EVERY
 410
     ! RECORD STRORES ONE MAGNITU
 420
     DE AND ONE PHASE DATA
 430
440 ASSIGN# 1 TO H$
450 DIM B8(144), B9(144), D8(144),
     D971441 | 144510=INTY2*K0*80
     +30
460 FR=01
470 FOR I=1 TO 800
480 DISP "FREQ LOOP=", I
490 F0=F0+Q4
500 DISP "FREQ (GHZ)=",F0
510 K1=.3/F0 ! WAVELENGTH
520 K0=X9/K1 | WAVE NUMBER
530 COSHR 650
540 DISP "E =",E0
550 DISP "P =",P0
560 PRINT# 1,1 ; E0,P0
570 NEXT I
580 ASSIGN# 1 TO *
590 CLEAR
600 DISP "END OF COMPUTATION"
610 END
620
630
640 L0=INT(2*K0*80+3)
650 IF L0K145 THEN 670
660 DISP "K0*A0 TOO LARGE FOR
    CURRENT ARRAY DIM"
    Z=K0*A0
679
680 GOSUB 910
690 F8=0
700 E9=0
710 FOR N=1 TO L0
720 L=L0-N+1
730 M8=08(L)^2+09(L)^2
740 M9=88(L)^2+89(L)^2
750 87=(L+.5)/M8/M9
768 A8=A7*(B9(L)*B9(L)-B8(L)*D8(
    LOG
778 89=87*(B8(L)*B9(L)+B9(L)*D8(
    LDD
780 FS=88-FS
798 F9=89-F9
800 NEXT N
810 F8=-F8
820 E9=-E3
830 E0=E8^2+E9^2
840 P0=ATN2(E9,E8)
850 E0=E0/K0%2%K1 ! CROSS-SECTIO
    М
860 P0=P0-X9#INT(P0/X9)
870 IF POKPI THEN 890
880 P0=P0-Y9
```

```
990 P0=-P0
900 PETHEN
910
920 IF Z>L0-1 THEN 1230
930 Z2=Z^2/2
940 N2=2*72+L0+1
950 D1=2*N2+3
960 D2=D1*(2*N2+5)
978 D3=D2#(2#N2+7)
980 04=03*(2*N2+9)
990 F1=1-Z2/D1+Z2^2/(2*D2)-Z2^3/
     (6*D3)
1000 F2=Z*(1/D1-Z2/D2+Z2^2/(2*D3
      )-Z2^3/(6*D4))
1010 M=2*72
1020 S1=F1
1030 F1=(2*M+1)*F1/Z-F2
1040 F2=S1
1050 IF ABS(F1)<1.E100 THEN 1090
1060 F1=F1*1.E-100
1070 F2=F2*1 E-100
1080 S1=S1*1.E-100
1090 M=M-1
1100 IF M+1>L0 THEN 1020
1110 B8(L0)=F2
1120 B8(L0-1)=F1
1130 NO=L0-2
1140 FOR K=1 TO NO
1150 N=L0-K-1
1160 B8(N)=(2*N+3)*B8(N+1)/Z-B8(
     N+20
1170 NEXT K
1180 A1=(SIN(Z)/Z-COS(Z))/B8(1)
1190 FOR K=1 TO L0
1200 BS(K)=A(*BS(K)
1210 NEXT K
     GOTO 1280
1220
1230 B8(1)=SIN(Z)/Z-COS(Z)
1240 B8(2) = (3/7/2-1) \pm SIN(7) - 3 \pm 0.0
     S(Z)/
1250 FOR N=3 TO LO
1260 B8(N)=(2*N-1)*B8(N-1)/Z-B8(
     N-25
1270 NEXT N
1280 B9(1)=-SIN(Z)-C0S(Z)/Z
1290 B9(2)=(1-3/Z^2)*C0S(Z)-3*SI
     NCZDZZ
1300 FOR N=3 TO L0
1310 B9(N)=(2*N-1)*B9(N-1)/Z-B9(
     N-20
1320 NEXT N
1330 B8(1)=(1-1/2^2)\pmSIN(2)+COS(
1340 D9(1)=(1/Z^2-1)*COS(Z)+SIN(
1350 FOR N=2 TO 19
1360 D8(N)=B8(N-1)-N*B8(N)/Z
```

1370 D9(N)=89(N-1)-N\*89(N)/Z 1380 NEXT N 1390 RETURN

```
10 | "CALIE DEIVEO"
 žā.
  76
    I CALIBRATION USING A SPHERE
 40 ! OVER L3-U9 GHZ AT F9 GHZ
      STEPS BASED ON THEORETICAL
 50
      VALUES COMPUTED USING
                             THE P
    ROGRAM "SPHERE DRIVEO"
 60
    I THE RESULTED SYSTEM TRANS-
      FER FUNCTION IS STORED AS:
 79
 80
       "CALIB3.DRIVE1" (3.187")
 96
100
    ! A$ IS THE FILE STORING THE
119
    I BACKGROUND DATA.
129
130
      C$ IS THE FILE STORING THE
    1
     SYSTEM TRANSFER FUNCTION
140
    1
    ! H# IS THE FILE STORING
150
    THEORETICAL DATA OF THE SPHE
    RE
169
170
      S# DESCRIBES THE SPHERE
189
199
200 C≉="CALIBZ DRIVE1"
210 H#="THEOP3 DRIVE1"
220 S#="3.187 INCH SPHERE"
230 A≰="BKGRND DRIVE1"
240
250 X9=2≮PI ! A PARAMETER
260 !
270 OPTION BASE 1
280 NO=3 ! NUMBER OF READINGS
290 | TAKEN AND AVERAGED FOR ONE
300 | FREQ.
310
320 F9=.1 ! FREQ. STEP IN GHZ
330
340 M1=51 / M1=(H9-L9)/E9+2
350 | NUMBER OF FREQ. CHECKED.
360 1
370 L9=10.1 ! LOWER FREQ.IN GHZ
380 U9=15 | UPPER FRED IN GHZ
398 1
400 DIM A(51.2) / BACKGROUND
                        DATA
410 DIM B(51.2) ! TARGET DATA
420 DIM G3(51/2) ! THEORY
    ! CPEATE C$.M1.16
470
440 ! CREATE A$, M1, 16
450
    ! STORE CALIBRATION AND BACK
    GROUND DATS IN A FILE OF MI
    RECORDS
    ! EACH RECORD CONTAINS ONE
460
    MAGNITUDE AND ONE PHASE
470 ! DATA AT A FREQUENCY
```

```
480
430 | READING THE THEORETICAL DA
     TS.
560
510 ASSIGN# 1 TO H$
520 K0=(L9-2-F9*2)*50
530 FOR T=1 TO M1
540 KA=KA+50*E9
550 READ# 1.K0 ; G3(I,1),G3(I,2)
560 NEXT I
570 ASSIGN# 1 TO *
580 GOSUB 2230 ! HEADER.
590 DISP "DO YOU WANT TO USE THE
     MOST RECENT BACKGROUND DATA
                         YZN "
600 INPUT P#
610 IF P$="H" THEN 690
626
639
    ! READING BACKGROUND DATA.
640
650 ASSIGN# 4 TO A$
660 READ# 4 ; A(;)
670 ASSIGN# 4 TO #
689 GOTO 949
690 CLEAR
700 REMOTE 7 ! REMOTE ALL
                    DEVICES
710 CLEAR 7 ! CLEAR ALL DEVICES
720 | INITIALIZE SIG.GEN TO FIRS
    T FREQ
239
    OUTPUT 719 / "P"/L9/"Z1K0L3M0
    N601"
740 CLEAR
750 DISP "REMOVE TARGET FROM CHA
    MBER, PUSH 'CONT' WHEN READY"
760 LOCAL 7
770 BEEP @ BEEP
780 PAUSE
790 REMOTE 7
800 CLEAR
810 DISP "TAKING BACKGROUND DATA
820 PRINT
830 | PRINT "BACKGROUND DATA"
840 PRINT
850 OUTPUT 719 ; "P", L9, "Z1K0L3M0
    NED1"
860 WAIT 200 ! WAIT FOR FREQ TO
     STABILIZE
870 GOSUB 1380
830
890 / STORING BACKGROUND DATA
900
910 ASSIGN# 7 TO AS
920 PRINT# 3 / A(,)
930 ASSIGN# 3 TO $
940 CLEAR
```

```
950 LOCAL 7
960 DISP "PUT TARGET INTO CHAMBE
     R. PUSH 'CONT' WHEN READY"
970 DISP "TARGET IS ",S$
980 PEEP @ BEEP
990 PAUSE
1000 REMOTE 7
1010 CLEAR
1020 DISP "COMPUTING TAPGET DATA
1939 PRINT
1848 | PRINT "TARGET DATA"
1050 PRINT
1060 OUTPUT 719 : "P", L9, "ZIK0L3M
      @N601"
1070 WAIT 200
1080 GOSUB 1970
1090 PRINT " "
1100 PRINT "TRANS, FUNCTION",S$
1110 PRINT " "
1120
1130 ! CALCULATE AND STORE TRANS
     FER FUNCTION
1149
1150 ASSIGN# 2 TO C$
1160 FOR M=! TO MI
1170 N1=8(M,1)-A(M,1)
1180 N2=8(M,2)-A(M,2)
1190 M6=G3(M, 1)/(N1^2+N2^2)
1200 X7=G3(M,2)-ATN2(N2,N1)
1210 X7=X7-X9*INT(X7/X9)
1230 15 X757 THEN X7=X7-X9
1230 PRINT# 2.M ; X6.X7
1240 ! PRINT USING 970 ; M.X6.X7
1250 IMAGE DD/1X/"X6="/SD.DDDE/1
     X."X7=".SD.DDDE
1260 NEXT M
1270 ASSIGN# 2 TO #
1280 CLEAR
1290 DISP "CALIBRATION COMPLETED
     DATA STORED IN",C$
1300 BEEP @ BEEP @ BEEP
1310 LOCAL 7
1320 END
1330
1340
1350
1360
1379
1380
     ! BACKGROUND DATA COLLECTIO
     N SUBROUTINE
1390
1400
       OUTPUT(L9-F9)TO U9 GHZ AT
      F9 GHZ STEPS
1410
     J=10*(L9-2*F9) ! FREQUENCY
     STARS AT L9-F9 GHZ
```

```
1420 FOR K=1 TO M1 ! NUMBER OF
     FREQUENCY STEPS
1430
     J=J+107F9
1440 IMAGE 19.32,148
1450 OUTPUT 719 USING 1440 / "P"
     J,"@@ZIK@L3M@N601"
1460 | TAKE DATA IN EROM 722%720
1470 GOSUB 1640
1480
1490 ! CALCULATE REAL&IMAGINARY
     FROM AMP. & PHASE
1500
1510 R1=A1 #C0S(P1)
1520 I1=81*3IN(P1)
1530 A(K, 1)=81
1540 A(K,2)=I1
1550
    ! PRINT USING 2040 ; A(K,1)
     A(K, 2)
1560
    NEXT K
1570 OUTPUT 719 : "P", L9, "Z1K0L3M
     9N601"
1580 RETURN
1590
1600
1618
1620
1630
     ! SUBROUTINE TO ENTER AMPLI
1640
     TUDE AND PHASE DATA FROM DI
     GITAL VOLIMETERS
1650
1660 ! PREPARE DIGITAL VOLTMETER
      TO SEND AMPLITUDE DATA
1670 ! NO READINGS TAKEN AND AVE
     RAGED FOR ONE EREG
1688
1690
1700 Visa | PARAMETERS FOR THE
1710 F1=0 ! AVERAGING PROCESS
1720 FOR L=1 TO NO
1730 OUTPUT 720 : "POFIBITIZIELOM
     ñ.
1740 WAIT 10
1750 ENTER 720 ; V
1760 WAIT 10
1770 OUTPUT 722 ; "F1R7T1M3A0H1"
1780 WAIT 10
1790 ENTER 722 ; F
1800 V1=V1+U
1810 F1=F1+F
1820 WAIT 10
1830 NEXT L
1840 V=V1/N0
1850 F=F1/N0
1860 A1=10^F | TRANSFER TO MAG
      FROM VOLTS.
```

```
1870 Pi=10010 / TRANSFER TO DEC
      FROM VOLTS.
1880 P1=DTR(P1)
1890 ! PRINT USING 1810 : K.At.P
1900 IMAGE DD, 2X, "A=", MD, DDDE, 2X
      "P=" SD DDDE
1918 PETHEN
1920
1939
1940
1950
1969
1970 ! TARGET DATA COLLECTION
       SUBPOUTINE
1980
1990 1
       OUTPUT(L9-F9)TO U9 GHZ AT
      F9 GHZ STEPS
2000
     J=10*(L9-2*F9) ! FREQUENCY
     STARS BT L9-E9 GHZ
2010 FOR K=1 TO M1 ! NUMBER OF
     FREQUENCY STEPS
2020 J=J+10xF9
2030 IMAGE 1A,3Z,14A
2040 OUTPUT 719 USING 2030 ; "P"
,J,"00Z1K0L3M0N601"
2050
    ! TAKE DATA IN FROM 722%720
     1820 GOSUE 1410
2060 GOSUB 1640
     ! CALCULATE REAL&IMAG FROM
2070
     AMP&PHASE
2080 R1=A1*C0S(P1)
2090 I1=A1*SIN(P1)
2100 B(K,1)=R1
2110 B(K,2)=I1
2120 ! PRINT USING 2040 ; B(K,1)
     /8(K,2)
IMAGE 4%, "R="/SD.DDDE/2%, "I
     =" .SD.DDDE
2140 NEXT K
2150 OUTPUT 719 ; "P", L9, "Z1K0L3M
     9N601"
2160 RETURN
2170
2180
2190
2200
2210
     ! HEADER SUBROUTINE
2220
2230 PRINT "
2240 PRINT
2250 CLEAR
2260 DISP "CALIBRATION STANDARD"
2270 PRINT "CALIBRATION STANDARD
     ",S#
```

```
10 | "TARGET DRIVEG"
  29 !
 30 1
      TARGET BACK-SCATTERING
 40 | USING C# DATA & STORE
      RESULTS IN G#
    + FREQUENCIES L9-U9 GHZ AT
 59
     F9
         GHZ STEPS
 60
 žā.
    ! FILE C# STORES THE SYSTEM
     TRANSFER FUNCTION
 89
 90 | FILE G# STORES TARGET DATA
     OBTAINED FROM THIS PROGRAM
1 99
110 ! FILE H& STORES THEORETICAL
     VALUES FOR PLOTTING OVERLAY
120
130 ! FILE A∜ STORES BACKGROUND
     DATA
149 1
150 C$="CALIBZ DRIVE!"
160 G#="SCR3 DRIVE1"
170 H#="THEOR3.DRIVE1:
180 A$="BKGRND DRIVE1"
199
200 ! CREATE G$,52,24
210 ! STORE TARGET DATA IN FILE
220 ! OF M1+1 RECORDS FIRST ONE
230 ! FOR THE AVERAGE PROCEDURE
      AND THE REST CONTAINS THE
249 I
250 | FREQUENCY MAGNETUDE AND
260 ! PHASE SHIFT.
270
280 ! CREATE A$.51,16
290 ! STORE CALIB. AND BACKGROUN
    DATA IN A FILE OF MI RECORDS
300 ! EACH RECORD CONTAINS ONE M
    AC AND PHASE AT A EREO.
319
320 OPTION BASE 1
330 NO=2 ! NUMBER OF READINGS
340 ! TAKEN AND AVERAGED FOR ONE
       FREQUENCY
359
360 DIM A(51.2) ! BACKGROUND DAT
370 DIM B(51.2) ! TARGET DATA
380 DIM G4(51-2) ! CALIBRATION
390 DIM N(51,3) ! RESULTANT
400 DIM M9(51,3)
410
420 N1=51
430 ! H1=(U9-L9):F9+2 NUMBER OF
440 ! FREQ. CHECKED
450 F9=.1 ! FREQ.STEPS IN GHZ.
```

```
460 09=15 | UPFER FREQ IN GHZ
470 L9=10 1 ! LOWER FREQ. IN GHZ
480 DIM T(800.2) | STORES THEORE
     TICAL DATA
490 X9=2*PI
500
510 / READING TRANSFER FUNCTION
520 !
530 ASSIGN# 1 TO C⊅
540 PEAD# 1
               64()
             - 1
550 ASSIGN# 1 TO #
560 ! MAT PRINT USING 330 : G4
570 IMAGE 2X, 3D, 40
580 1
590 REMOTE 7 ! REMOTE ALL
                      DEVICES
600 CLEAR 7 ! CLEAP ALL DEVICES
610 OUTPUT 719 ; "P1Z1K0L3M0N601"
       INITIAL SETUP OF 719
620 CLEAR
630 DISP "DO YOU WANT TO USE THE
     MOST RECENT BACKGROUND DATA
            YZN"
649 INPUT P$
650
560 IF P#="N" THEN 740
670
580 ! READING BACKGROUND DATA
690 i
700 ASSIGN# 4 TO AM
710 PEAD# 4 / 8(/)
720 ASSIGN# 4 TO #
730 GOTO 900
740 DISP "REMOVE TARGET FROM
    CHAMBER, PUSH 'CONT' WHEN REA
    DY"
750 LOCAL 7
760 BEEP @ BEEP
770 PAUSE
780 DISP "TAKING BACKGROUND DATA
790 REMOTE 7
800 OUTPUT 719 ;"P"/L3,"Z1K9L3M9
    N601" ! INITIAL SETUP OF 719
SIG WAIT 199
820 GOSUB 2560
830
840 / STORING BACKGROUND DATA.
850
860 ASSIGN# 5 TO A≸
870 PRINT# 5 : 8(.)
880 ASSIGN# 5 TO *
890 CLEAR
900 DISP "PUT TARGET INTO CHAMBE
     F. PUSH 'CONT' WHEN READY"
910 LOCAL 7
920 BEEP @ BEEP
930 PAUSE
```

```
946 REMOTE 7
950 OUTPUT 719 ;"P",L9,"Z1K0L3M0
N601" ; INITIAL SETUP OF 719
960 WAIT 500
978 GOSUB 3418
980 CLEAR
990 DISP "COMPUTING TARGET DATA"
1000 GOSUB 3130
1010 1
1020 !
       COMPUTING TARGET DATA
1030 | WITHOUT BACKGROUND AND TH
1949
     I FREG FOR EACH RECORD.
1050 F0=L9-2#F9
1060 FOR M=1 TO N1
1070 F0=F0+F9
1080 N(M, 1)=F0
     87=8(M,1)-A(M,1)
1090
1100 X8=B(M.2)-A(M.2)
1110 X6=(X7^2+X8^2)*G4(M,1)
1120 N(M,2)=X6
1130 X8=ATN2(X8,X7)+G4(M,2)
     X8=X8-X9*INT(X8/X9)
1140
1150 IF X8>PI THEN X8=X8-X9
1160 N(M,3)=28
1170 NEXT M
1180 DISP "PRINT DATA? Y/N"
1190 BEEP @ BEEP
1200 INPUT PE
1210 IF P#="H" THEN 1250
1220 PRINT "
                   FREG
                               CRSEC
          PHASE"
1230 MAT PRINT USING 1240 : N
1240 IMAGE 2X, 3D, 4D
1250 CLEAR
1260 LOCAL 7
1270 DISP "PLOT MAGNITUDE FOR
           THIS MEASUPMENT? YAN"
1280 INPUT PS
1290 IF P#="N" THEN 1330
1300 DISP "SELECT PEN. PUSH
'CONT' WHEN READY"
1310 PAUSE
1320 GOSUB 3570
1330 GUEAR
1330 CLEAR
1340 DISP "PLOT PHASE FOR THIS
MERSUPMENT ? YZN"
1360 INPUT P#
1370 IE P#="N" THEN 1410
1380 DISP "SELECT PEN. PUSH
           'CONT' WHEN READY"
1390 PAUSE
1400 GOSUS
            4610
1410 CLEAR
```

```
1420 DISP "DO YOU WANT TO MAKE
           AVERAGE WITH PREVIUSE
     Detac"
1430 DISP "2Y-N"
1440 BEER & BEER
1450 INPUT P#
1460 IF P#="Y" THEN 1750
1470 DISP "DO YOU WANT TO STORE
          DATA 2 YAN "
1480 INPUT P#
1490 IF P#="N" THEN 2280
1500 MA=1
1519 DISP "DO YOU WANT TO STORE
           DATA IN FILE"
1528 DISP
           Gŧ
1530 DISP "? Y-N"
1540 INPUT P#
1550 IF P$="Y" THEN 1626
1560 DISP "ENTER NAME OF THE DAT
     A FILE TO BE USED FOR STORE
     GE"
1570 INPUT GE
1580 DISP "IS THIS AN OLD FILE
TO BE UPDATED ? Y/N "
1590
    INPUT PS
1600 IF P#=""" THEN 1620
1610 CREATE G#, 53, 24
1620 DISP "ENTER LENGTH OF TARGE
1630 BEER @ BEER
1649 INPUT M1
1650 DISP "ENTER DIAMETER OF TAR
     GET"
1660 BEEP @ BEEP
1670 INPUT M2
1689 (
1690 ! STORE MEASURED DATA
1700 1
1710 ASSIGN# 2 TO G$
1720 PRINT# 2 ; M0,M1,M2,N(,)
1730 ASSIGN# 2 TO *
1740 GOTO 2280
1750 DISP "DOES THE DATA STORED
           IN FILE"
1760 DISP G#
1770 DISP "? Y/N"
1780 INPUT P#
1790 IF P#="Y" THEN 1870
1300 DISP "ENTER NAME OF DATA
          FILE TO BE USED FOR THE
           AVERAGE"
1810 !
1820 INPUT G≴
1839
1340 ! READ OLD DATA
1850 ! AND MAKES WIGHTED AVERAGE
1860 ! WITH NEW DATA.
```

```
1870 ASSIGN# 6 TO G#
1880 READ# 6 : M0.M1.M2.M9(,)
1890 ASSIGN# 6 TO *
1900 FOP K=1 TO NI
1910 M9(k)2)=M9(K)2)*M0+N(K,2)
1920 M9(K,2)=M9(K,2)/(M0+1)
1930 M9(K,3)=M9(K,3)*M0+N(K,3)
1940 M9(K,3)=M9(K,3)/(M0+1)
1950 N(K/2)=M9(K/2)
1960 N(K,3)=M9(K,3)
1970 NEXT K
1980 M0=M0+1
1990
2000 ! STORE NEW AVERAGE
2010 ASSIGN# 7 TO G≸
2020 PRINT# 7 ; M0,M1,M2,N(,)
2030 ASSIGN# 7 TO *
2040 PRINT "DATA IS AVERAGE OF",
     MO. "MEASHPMENTS"
2050 DISP "PRINT DATA? Y/N"
2060 BEEP @ BEEF
SAZA INPUT P#
2080 IF P#="N" THEN 2126
     PRINT "
                  FREQ
                             CRSEC
2090
          PHASE"
2100 MAT PRINT USING 1240 ; N
2110 IMAGE 2X,3D,40
2120 DISP "PLOT MAGNITUDE?
                             Y / N "
2130 BEEP @ BEEP
2140 INPUT PF
     IF P$="N" THEN 2190
2150
2160 DISP "SELECT PEN FOR MAGNIT
     HOE PLOT
                PUSH 'CONT' WHEN
     READY
2170 PAUSE
2180 GOSUB 3570
2190 CLEAR
200 DISP "PLOT PHASE?
                         Y 2 N **
 210 BEEP @ BEEP
2220
    INPHIT PE
2239
    TE P#="N" THEN 2276
    DISP "SELECT PEN AND CHANGE
PAPER FOR PHASE PLOT. PUS
2240
     H 'CONT' WHEN READY."
2250 PAUSE
2260 GOSUB 4610
2270 CLEAR
2280 DISP "DO YOU WANT TO"
2290 DISP "OBTAIN DATA"
2300 DISP "FOR A NEW TARGET?"
2310 DISP
          11
    DISP "ENTER Y/N"
2329
2330
     THRUT P#
2340
     IE P#="N" THEN 2479
2350
     DISP "DO YOU WANT TO USE
      THE SAME FILE"
2360 DISP G#
```

```
2370 DISP "TO STORE NEW DATAS Y"
     Ha
2380 INPUT PE
2390 IF P#="Y" THEN 2460
2400 DISP "ENTER NEW FILE NAME T
     O STORE TARGET DATA"
     INPUT G#
2410
2420 DISP "IS THIS AN OLD FILE
     TO BE UPDATED? YZN"
2430
     INPUT PE
2440 IF P#="Y" THEN 2460
2450 CREATE G$.52.24
2460 GOTO 598
2470 CLEAR
2480 DISP "END OF PROGRAM"
2490 BEER @ BEER @ BEER
2500 END
2518
2520
2530
2540
2550
2560
       BACKGROUND DATA COLLECTIO
     N SUBPOUTINE
     I OUTPUT(L9-F9) TO U9 GHZ
     J=10*(L9-2*F9) ! FREQUENCY
2580
     STARTS AT L9-F9 GHZ TO BE
     INCREASED AT F9 GHZ STEPS
2590 FOR K=1 TO N1 ! NUMBER OF F
     REQUENCY STEPS
2600 J=J+10#F9
2610 IMAGE 1A.32,14A
2620 OUTPUT 719 USING 2610 ; "P"
     JJ, "00Z1K0L3M0N6O1"
2639
     ! 50 MBEC WAIT FOR FREQUENC
     Y TO STABILIZE
2640 WAIT 50
2659
     ! TAKE DATA IN FROM 722 AND
      729
2660
    G0SUB 2820
2670
     ! PEAL AND IMAGINARY PARTS
2688
       FROM AMP. AND PHASE
2690 R1=A1 C08(P1)
2700 I1=81#SIN(P1)
2710 A(K,1)=R1
2720 A(K,2)=I1
     ! PRINT "I1=",A(K,2)
2730
     ! PRINT "R1="JA(K)1)
2740
2750 NEXT K
2760 OUTPUT 719 ; "P", L9, "Z1K0L3M
     0N601" ! INITIAL SETUP OF 7
     19
2770 RETURN
2780
2790 I
```

```
2800
2810
2820
     I SUPPOSITIVE TO ENTER AMPLI
     TUDE AND PHASE DATA FROM DI
     GITAL VOLIMETER
2830
2840
     ! PREPARE DIGITAL VOLTMETER
      TO SEND AMPLITUDE DATA
2850
     ! NO READINGS TAKEN AND AVE
     RAGED FOR ONE FREQUENCY
2860 V1=0 ! PARAMETERS FOR AVERA
                GING PROCESS.
2870 W1=0
2880 FOR L=1 TO NO
2890 OUTPUT 720 ."FIRITIZIFLOMO"
2900 WAIT 10
2910 ENTER 720 ; VO
2920 WAIT 10
2930 OUTPUT 722 : "F1R7T1M3A0H1"
2940 WAIT 10
2950 ENTER 722 ; WO
2960 V1=V1+V0
2970 W1=W1+W0
2980 WAIT 10
2990 NEXT L
3000 V0=V1/N0
3010 NO=N1/NO
3020 | TRANSFERS FROM VOLTS TO
          AMPL
3030 A1=10^W0
3040 | TRANSFERS TO DEG. FROM VO.
          LTS
ZOSO Pi=tookyo
3060 | PRINT "A1=".A1
3070 P1=DTR(P1)
3080 | PRINT "P1=".P1
3090 RETURN
3100
3110
3120
3130 ! DATA COLLECTION SUBROUTIN
     ! OUTPUT(L9-F9)TO U9 GHZ AT
3140
      F9 GHZ STEPS
3150 J=10*(L9-2*F9) ! INITIAL FR
     EQUENCY AT L9-F9 GHZ
3160 FOR K=1 TO N1 ! FREQUENCY S.
     TEPS
3170
     DE LETAKER L ER GHZ INCREME
     NTS
3180 IMAGE 18,3Z,148
3190 OUTPUT 719 USING 3180 ; "P"
     J, "00Z1K0L3M0N6O1"
3200
     ! 50 MSEC WAIT FOR FREQUENC
     Y TO STABILIZE
3210 WALT 50
```

```
3220 | TAKE DATA IN EROM 7224720
3230 GOSHB 2320
7740
3250
     ! REALGIMAG FROM AMP %PHASE
3260 R1=81%C09(P1)
327B
     T1=AlkSIN(P1)
3280
     B(K,1)=R1
3290 B(K,2)=I1
     | PRINT "R1=",B(K,1)
3300
3310
     ! PRINT "I1=",B(K,2)
3320 NEXT K
     OUTPUT 719 : "P",L9, "Z1K0L3M
0N601" ! INITIAL SETUP OF 7
3330
      19
3340 PETURN
3350
3360
รีรีร์ดิ
3380
     ! HEADER SUBROUTINE
3390
3400 D$="MONTH/DATE/YEAR"
3410 PRINT "
3420 PRINT " "
3430 CLEAR
3440 DISP "ENTER TODAY'S DATE -
     MONTH, DATE, YEAR"
3450 INPUT D#
3460 DISP "ENTER TGT DESCRIPTION
3470 INPUT T#
3480 PRINT
            Đέ
3490 PRINT
            "TARGET IS ".TS
3500 PRINT
            3510 PRINT "+ k**************
3520 PRINT "
3530 CLEAR
3540 RETURN
3550
3560
3570
     ! MAGNITUDE PLOTTING
3580
     I SUBROUTINE
3590
3600 PLOTTER IS 705
3610 LOCATE 32,122,20,85
3620 FRAME
3639
       SEARCH FOR MAX & MIN
3649
       S0=N(2,2)
3650
     1 81=80
     ! FOR M=3 TO N1
3660
     ! IF $00 N(M, 2) THEN $0=N(M,
3670
3680
     ! IF $1/N(M,2) THEN $1=N(M.
     Ź١
3690 ! NEXT N
3700 DISP "ENTER LOWER VALUE FOR
           MAGNITUDE PLOTTING
```

```
3710 BEER @ BEER
3720 IMPUT SA
3730 DISP "ENTER HERER VALUE FOR
           MAGNITUDE PLOTTING"
3740 INPUT SI
3750 L1=INT(L9)
3760 U1=CEIL/U9)
3779
     ! CALCULATE SCALE STEPS
3780
3790
     ! FOR MAGNETUDE
3800 S3=LGT(S1)
     S4=INT(S3)-1
3810
3820
     85=93-94
3830 S5=INT(10^S5)+1
3940 94±10094
3850 L0=INT(S0/S4)
3860 IF S5-L0<=14 THEN 3940
     IF S5-L0>=50 THEN 3910
3870
3888 55= 5 $85
3890 $4=$4*2
3900 GOTO 3850
3910 85= 2#85
3920 $4=5#84
3930 6010 3850
3940 L0=S4#L0
3950 10=95#34
3960 D0=U0-L0
3970 SCALE L1,U1,L0,U0
3980 FXD 0,4
3990 LAXES -1,54,L1,L0
4000 MOVE L1,0
4010 FOR K=2 TO N1
4020 N9=N(K,1)
4030 R9=N(K)20
4040 GOSUB 4510
4050 NEXT K
4060 M5=(U1+L1)/2
4070 MOVE M5, L0-.09*D0
4080 LORG 5 @ CSIZE 3.1.0
4090 ! LABEL "FREQUENCY (GHZ)"
4100 MOVE L1-1..5*(L0+00)
4110 LDIR PI/2
     ! LABEL "CROSS SECTION (SQ. METERS)"
4120
     MOVE M5, U0+ 09*D0
4130
4140 LDIR 0
4150 CSIZE 3,1,0
4160 LABEL T#
4170 MOVE M5. U0+.03*D0
4180 LABEL D#
4190 PENUP
4200 DISP "OVERLAY THEORETICAL C
     URVE? YZN"
4210
4220 INPUT P#
4230 IF P#="N" THEN 4496
```

```
4240 DISP "IS THE THEORETICAL
DATA STORED IN THE FILE".H$
4250 DISP "7 VAN"
4260 INPUT P#
4270 IF P#="Y" THEN 4300
4280 DISP "ENTER NAME OF THE
     DATA FILE TO BE PLOTTED. "
4290 INPUT H$
4300 BEEP @ BEEP
4310 DISP "CHANGE PEN IF DESIRED
        PUSH 'CONT' WHEN PERGY "
4320 PAUSE
4330 89916N# 3 TO H$
4349 .11=(19-2)*50
4350 32=(09-2) $50
4360 FOR J=JI TO J2
4370 READ# 3.J ; T(J,1),T(J,2)
4380 NEXT J
4390 ASSIGN# 7 TO #
4400 F0=L9
4410 R9=T(J1,1)
4420 MOVE F0.R9
4430 FOR I=J1+1 TO J2
4440 F0=F0+.02
4450 R9=T(1.1)
4460 DRAW FOLES
4470 NEXT I
4480 PENUP
4490 PETHEN
4500
4510 ! PLOT CROSS
4520 MOVE M9 P9
4530 CSIZE 2, 5,0
4550 ! IMOVE
              00025,.00025
       IDRAN - .0005.0
4560 L
4570 RETURN
4580
4590
4600
4610 ! PHASE PLOTTING SUBROUTINE
4620
4630 PLOTTER IS 705
4640 LOCATE 32,122,20,85
4650 FRAME
4660 84=.25
4670 U0=1
4680 L0=-1
4690 D0=U0-L0
4700 SCALE L1, U1, L0, U0
4710 FXD 0.3
4720 LAXES -1,84,L1,L0
4730 MOVE L1/0
4740 FOR K=2 TO N1
4750 M9=N(K,1)
4760 R9=N(K, 3)/PI
4770 GOSUB 5110
```

```
4780 NEXT K
4790 MÖVE M5.L0- 09≭00
4800 LORG 5 @ CSIZE 3.1.0
     ! LABEL "FREQUENCY (GHZ)"
4810
4820 MOVE L1-1, .5*(L0+U0)
4830 LDIR PI - 2
4840 LABEL "PHASE (PI)"
4850 MOVE M5. U0+ 09*D0
4860 LDIR 0
4870 CSIZE 3,1,0
4880 LABEL T#
4890 MOVE M5, U0+.03*D0
4900 LABEL D≢
4910 PENUP
4920 DISP "OVERLAY THEORETICAL C
     HRUE? YZN"
4930
     INPUT P#
4940 IF P#="N" THEN 5070
4950 BEEP @ BEEP
     DISP "CHANGE PEN IF DESIRED
4960
        PUSH 'CONT' WHEN READY."
4970 PAUSE
4980 FR=19
4990 R9=T(J1,2)/PI
5000 MOVE F0,R9
5010 FOR I=J1+1 TO J2
5020 F0=F0+.02
5030 R9=T(I,2)/PI
5040 DRAW F0, R9
5050 NEXT I
5060 PENUP
5070 RETURN
5080
5090
5100
5110 ! PLOT DOT
5120 MOVE M9.R9
5130 CSIZE 2, 5.0
5140 LABEL "*"
5150 ! IMOVE .00025,.00025
5160 ! IDRAW -.0005.0
5170 RETURN
```

MECHANICAL SPECIFICATIONS:

-	Absorber	Absor	Absorber Height (In.)	(In.)	Pyramid Base	Pyramids
8-1/2 1-1/2 7	Size (In.)	Overall	Base	Pyramid	Size (In.)	Per Absorber
	24 × 24	8-1/2	1-1/2	7	3×3	64

MAXIMUM REFLECTION AT NORMAL INCIDENCE:

Ku X C Band Band Band	S	1	900	000		
Band	0			2005	200	120
Management and Association of the Control of the Co	Duna	Band	MHz	MHz	MHz	MHz
50 db 50 db 45 db	40 db	30 db				

Specifications of the Absorber Material of the Floor

MECHANICAL SPECIFICATIONS:

Pyrainid Base Pyramids	mid Size (In.) Per Absorber	5 6×6 16
Absorber Height (In.)	Base Pyran	2-1/4 16
Absor	Overall	18-1/4
Absorber	Size (In.)	24×24

MAXIMUM REFLECTION AT NORMAL INCIDITICE:

0	2	
120	HW	
200	MHz	
300	MHz	
500	MHz	30 db
1	Band	40 db
S	Band	45 db
С	Band	50 db
×	Band	50 db
Κυ	Band	db 05

Specifications of the Absorber Material of the Walls

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"Head-on" Scattering of a tubular cylinder of finite length for radar target identification purposes.

: 564

Thesis G25958 c.1

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"Head-on" Scattering of a tubular cylinder of finite length for radar target identification purposes.

